

## Compton Electrons in POE Range Telescope Scintillator

Probability of Gamma Ray passing through anti Compton Scintillator:

This scintillator is  $\frac{3}{4}$ " Plastic (NE102) upper, Anthracene with threshold of  $\frac{1}{3}$  Minimum Ionizing protons

Scintillator Data (From Nuclear Enterprises Ltd NE102 Plastic Fluorescence Bulletin # 1004)

$$\text{Density} = 1.03 \text{ gm/cc}$$

$$\text{Electrons per cc.} = 3.9 \times 10^{23}$$

Atomic composition

$$\text{H} - 5.25 \times 10^{22} \text{ atoms/cc}$$

$$\text{C} - 4.75 \times 10^{22} \text{ "}$$

$$\text{N} - 1.6 \times 10^{19} \text{ "}$$

$$\text{O} - 1.8 \times 10^{19} \text{ "}$$

Table 2 p 17 of NBS Circular 583 (X-Ray Atten. Coeffs from 10KeV to 10MeV) Give the cross section calculated from the Klein - Nishina formula for Compton interactions in cm<sup>2</sup>/electron, convert to cm<sup>2</sup>/g.

$$\text{ie } \sigma(\text{cm}^2/\text{g}) = \sigma(\text{cm}^2/\text{elec}) \cdot \frac{\text{Electrons}}{\text{cm}^3} \cdot \frac{1}{\text{Density}}$$

$$\text{The threshold of the scintillator is } \frac{1}{3} (1.015 \frac{\text{MeV}}{\text{g/cm}^2}) (0.98 \frac{\text{g}}{\text{cm}^3}) = \frac{10.9 \text{ MeV}}{3} = 3.63 \text{ MeV}$$

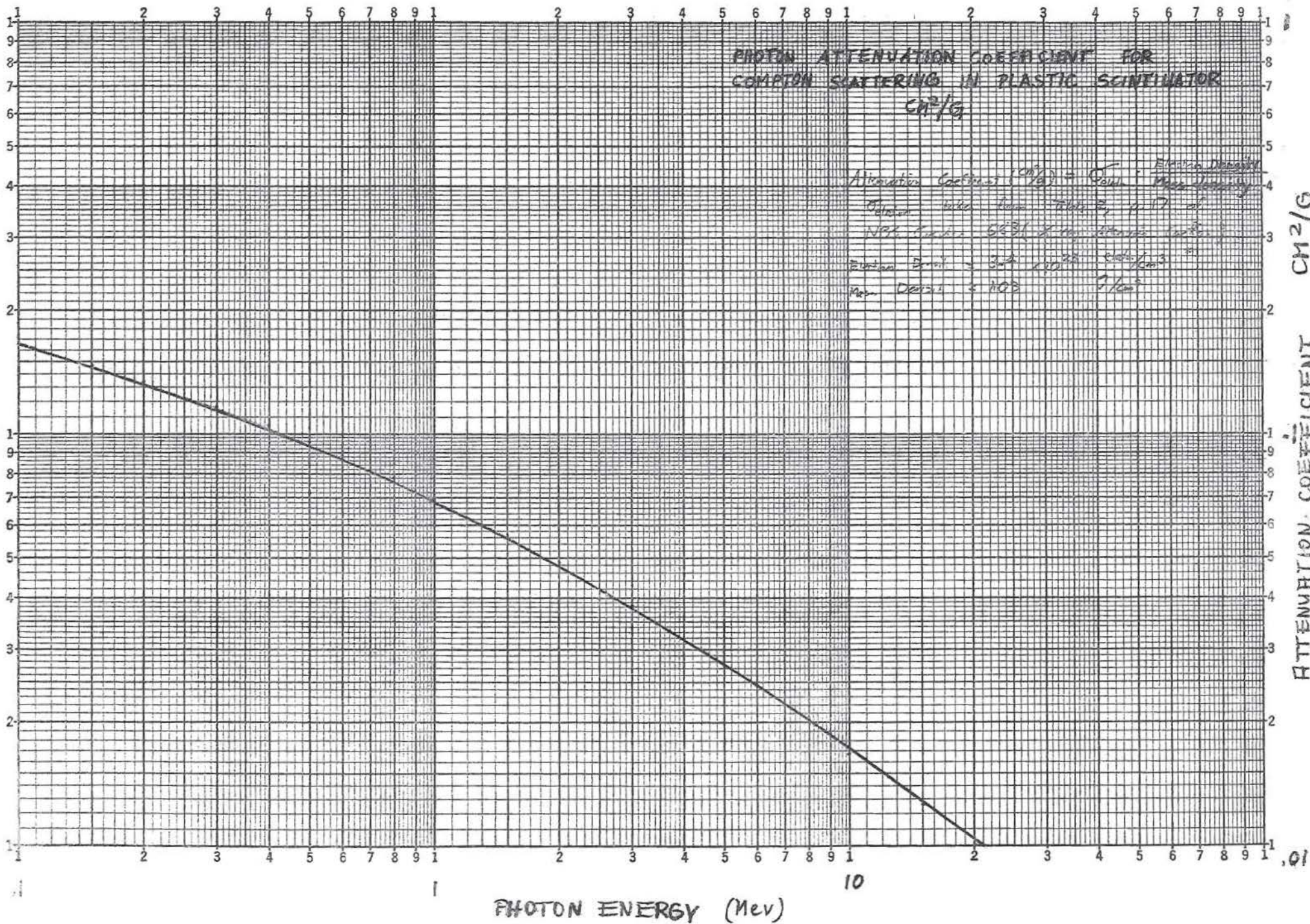
See Graph Next Page

Graph on next page shows that there is generally less than 10% chance of gamma interaction with scintillator.

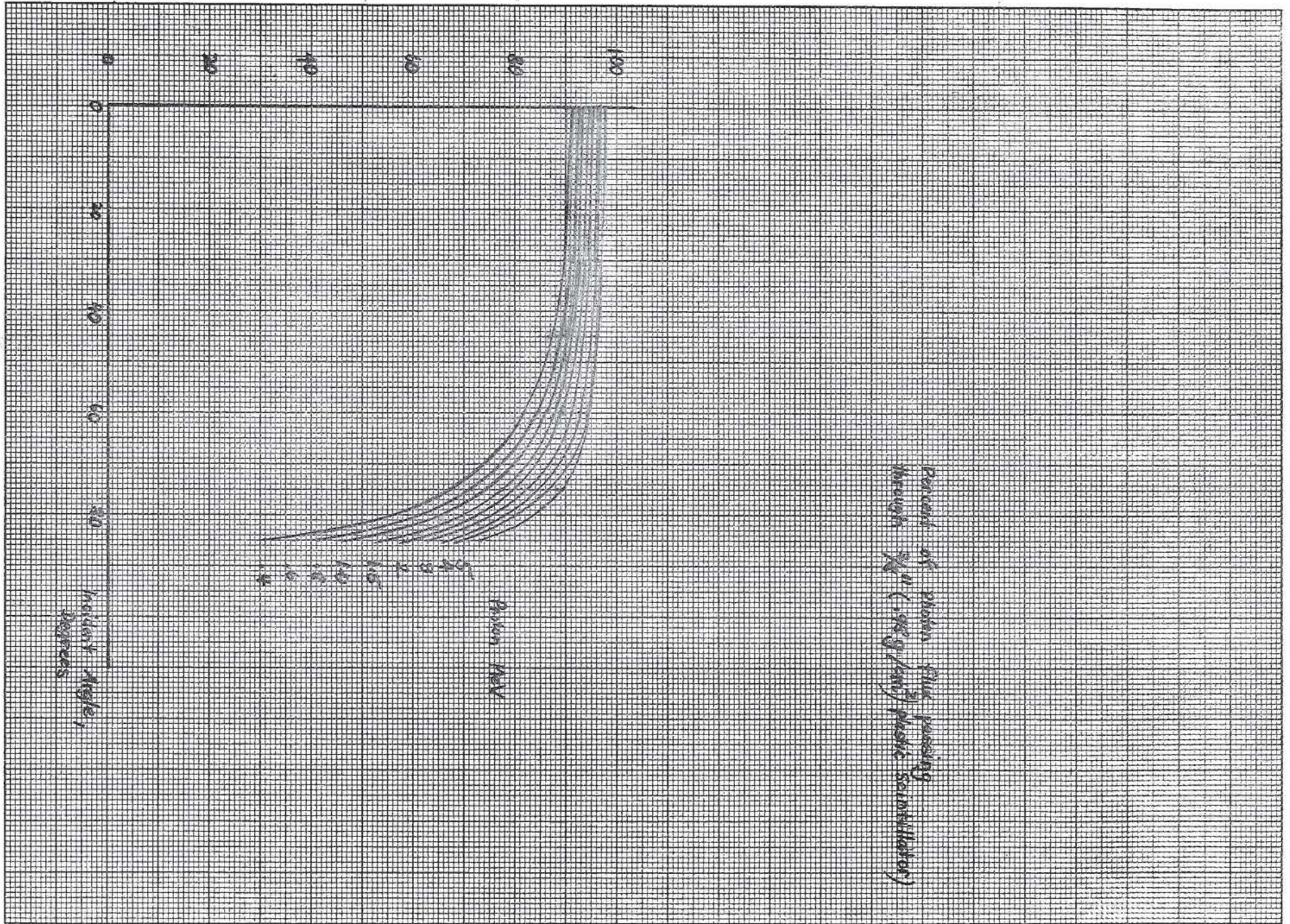


PHOTON ATTENUATION COEFFICIENT FOR  
 COMPTON SCATTERING IN PLASTIC SCINTILLATOR  
 CM<sup>2</sup>/G

Attenuation Coefficient (CM<sup>2</sup>/G) =  $\frac{\text{Electron Density}}{\text{Mass Density}}$   
 Electron Density =  $\frac{\text{Number of Electrons}}{\text{Volume}}$   
 Mass Density =  $\frac{\text{Mass}}{\text{Volume}}$   
 Plastic Scintillator: 56.3% C, 43.7% H  
 Electron Density =  $3.2 \times 10^{23} \text{ electrons/cm}^3$   
 Mass Density =  $1.03 \text{ g/cm}^3$









### Solid Angle

$$\text{Solid angle} = \Omega = 2\pi \int_0^\theta \sin\theta d\theta = 2\pi(1 - \cos\theta)$$

$\Omega/4\pi$	$\theta$
.05	25.8°
.1	36.9
.15	45.6
.2	53.1
.25	60.0
.3	66.4
.35	72.5
.4	78.5
.45	84.3
.5	90.0

Count rate spectrum for D2 single events Assuming

no photon attenuation from shielding or anticoincidence.  
 Procedure: Consider 5 conical zones such that each has solid angle of  $\frac{1}{10} \times 4\pi$ .

Zone	Small cone angle	Large cone angle	Median angle
1	0.0	36.9	25.8
2	36.9	53.1	45.6
3	53.1	66.4	60.0
4	66.4	78.5	72.5
5	78.5	90	89.3

The Count rate for each zone is determined as follows for a given energy interval

$$R = (\text{Flux} \times \text{Energy Interval}) \cdot (\text{Geometric Factor}) \cdot (\text{Interaction Probability})$$

$$\text{Geometric factor} = \text{Area} \times d\Omega \times \cos(\text{Median Angle})$$

$$\text{Interaction probability} = 1 - e^{-UX \sec(\text{Median angle})}$$

$$\text{Area detector area} = 4.5 \text{ cm}^2 \text{ for D2}$$

U = Attenuation Coefficient (see graph)

$$X = \text{detector thickness} = 2.42 \text{ g/cm}^2 \text{ for D2}$$

Counting rate through top window

Zone		1	2	3	4	5	
Solid Angle Increment		$4\pi/10$	$4\pi/10$	$4\pi/10$	$4\pi/10$	$4\pi/10$	
Median Angle		$28.5^\circ$	$45.6$	$60$	$72.5$	$84.3$	
Geometric factor		4.77	3.96	2.82	1.70	.565	
X. sec (Median Angle)		.275	.346	.485	.805	2.49	(X. sec)
Photoxy Energy	Atten. Coef	Interaction Probability = $1 - e^{-UX_{sec}(\text{Median Angle})}$					
	U	Zone 1	2	3	4	5	
.4 MW	.096	.027	.033	.043	.074	.209	
.6	.080	.022	.027	.034	.062	.173	
.8	.070	.019	.024	.031	.055	.157	
1.0	.064	.017	.022	.031	.051	.144	
1.3	.056	.015	.019	.027	.045	.126	
2	.045	.012	.016	.022	.036	.104	
3	.035	.010	.012	.017	.028	.082	
4	.029	.008	.010	.014	.023	.064	
5	.025	.007	.009	.012	.020	.056	
Photon Energy Interval	Flux $\text{sec}^{-1}$ Photons in Energy Interval	Counting rate $\text{sec}^{-1}$					@ $4\mu\text{Ci}$ attached
		1	2	3	4	5	Total
.4 $\pm$ .1	1.40	.187	.182	.169	.175	.165	.878
.6 $\pm$ .1	.44	.048	.047	.047	.046	.043	.231
.8 $\pm$ .1	.22	.021	.020	.021	.020	.019	.101
1.0 $\pm$ .1	.16	.013	.014	.014	.014	.013	.068
1.3 $\pm$ .2	.21	.013	.016	.016	.016	.015	.076
2 $\pm$ .5	.35	.021	.022	.022	.022	.020	.107
3 $\pm$ .5	.27	.013	.013	.013	.013	.012	.064
4 $\pm$ .5	.23	.009	.009	.009	.009	.009	.045
5 $\pm$ .5	.20	.007	.007	.007	.007	.007	.035
							1.705

Note that Compton interactions are almost negligible for this detector case

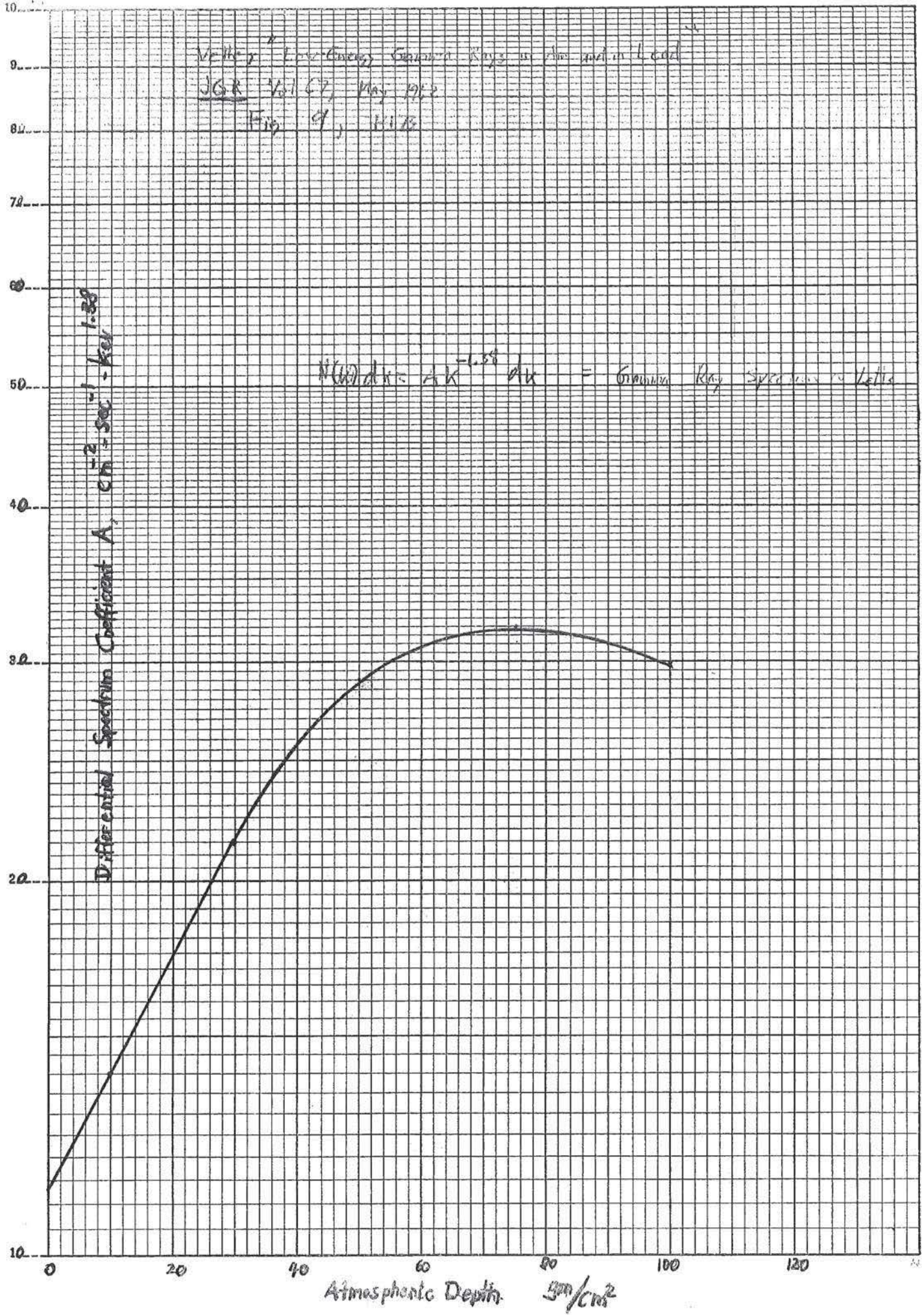


KE SEMI-LOGARITHMIC 46 4653  
1 CYCLE X 70 DIVISIONS MADE IN U.S.A.  
KEUFFEL & ESSER CO.

Vetter, "Low Energy Gamma Rays in Air and in Lead"  
JGR Vol 67, May 1952  
Fig 9, 1113

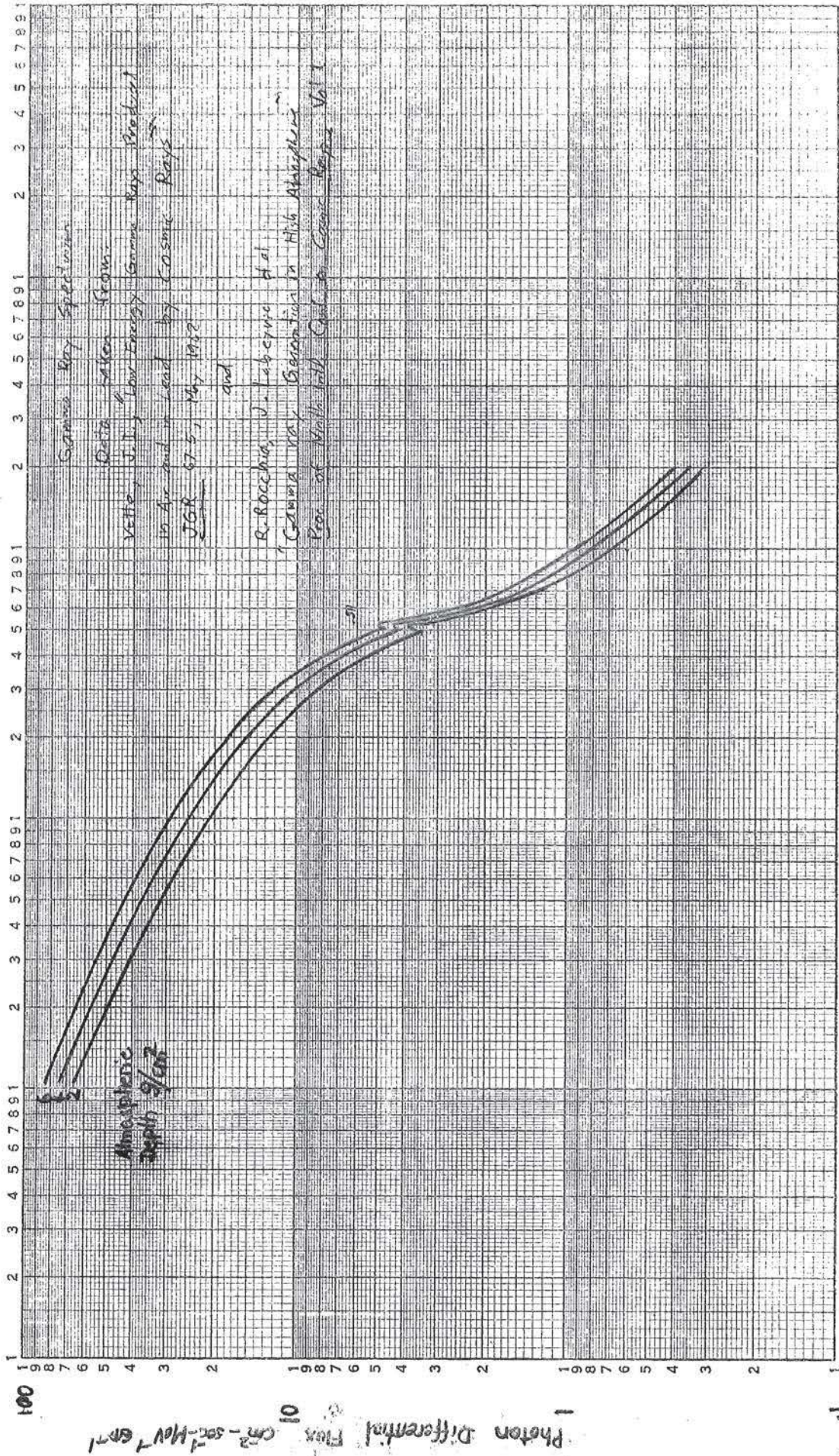
Differential Spectrum Coefficient  $A$ ,  $\text{cm}^{-2} \text{sec}^{-1} \text{kev}^{-1.30}$

$N(x) dx = A x^{-1.30} dx = \text{Gamma Ray Spectrum in Lead}$



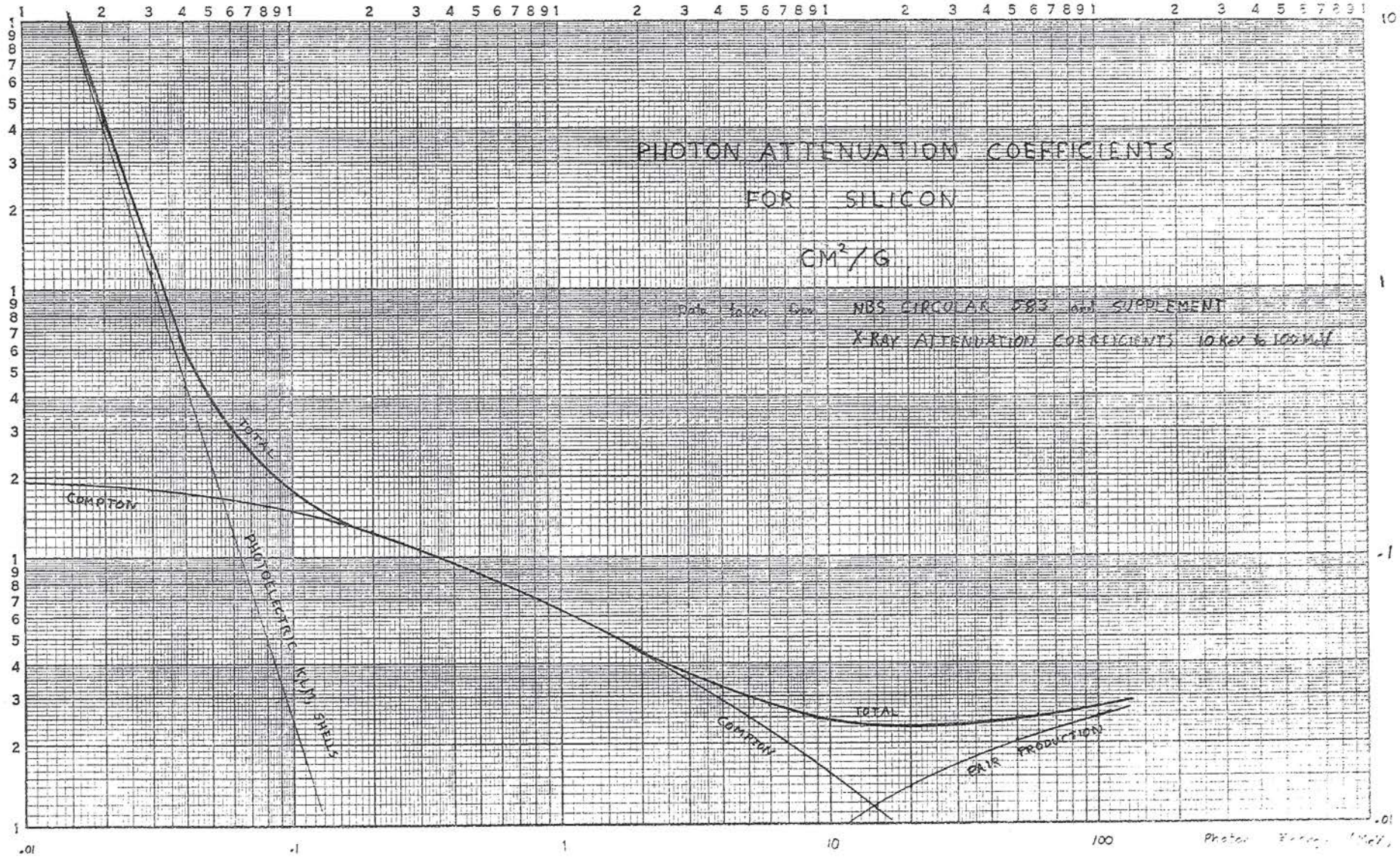
Atmospheric Depth  $\text{g}/\text{cm}^2$





10  
 1  
 .1  
 .01





Attenuation Coefficient,  $\text{cm}^2/\text{g}$

Photon Energy, MeV



Now find distribution of recoil detection times, by using the energy-angle probability graph

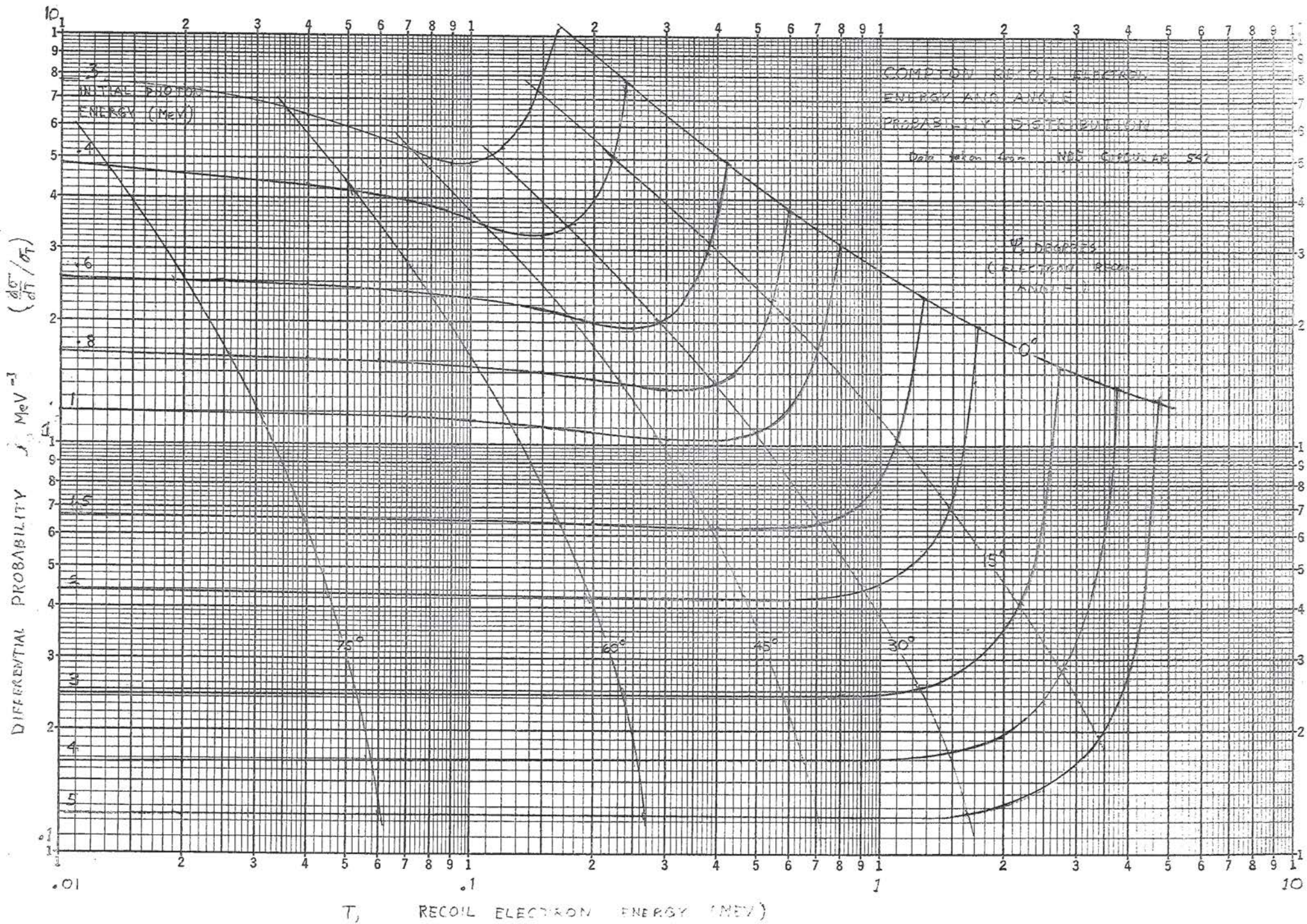
Electron Energy (MeV)	0-.5	.5-1	1-1.5	1.5-2	2-2.5	2.5-3	3-3.5	3.5-4	4-4.5	
Photon Energy	Counts Per Second in each interval									
.4 E <sub>1</sub>	.878	0	0	0	0	0	0	0	0	
.6 E <sub>1</sub>	.231	0	0	0	0	0	0	0	0	
.8 E <sub>1</sub>	.076	.025	0	0	0	0	0	0	0	
1 E <sub>1</sub>	.034	.034	0	0	0	0	0	0	0	
1.3 E <sub>1</sub>	.030	.040	.006	0	0	0	0	0	0	
2 E <sub>1</sub>	.023	.023	.030	.028	0	0	0	0	0	
3 E <sub>1</sub>	.008	.008	.008	.010	.015	.015	0	0	0	
4 E <sub>1</sub>	.004	.004	.004	.004	.005	.006	.009	.010	0	
5 E <sub>1</sub>	.002	.002	.002	.002	.002	.003	.003	.004	.004	
	1.286	.136	.050	.044	.022	.024	.012	.019	.015	

(This is for Top Hemisphere only)

Recoil Electron Energy spectrum for Compton incident neutrons over 4% straggling

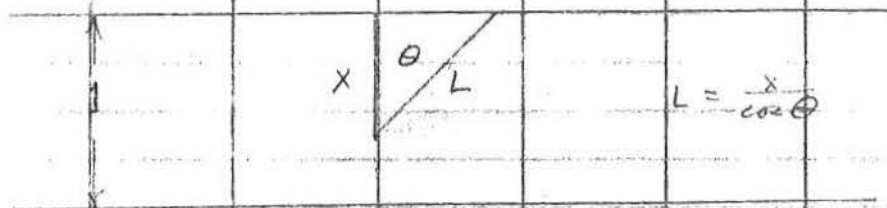
Recoil Electron Energy (MeV)	0-.5	.5-1	1-1.5	1.5-2	2-2.5	2.5-3	3-3.5	3.5-4	4-4.5
Counting Rate Mev <sup>-1</sup> Sec <sup>-1</sup>	5.05	.545	.200	.176	.088	.096	.047	.050	.000







Exit path length probability, for events occurring isotropically in a slab of unit thickness and infinite extent.



Assume  $L$  is constant, then the probability of an exit length  $L$  at a depth  $x$  is

$$P(L)_x dL = P(\theta) \frac{d\theta}{dL} dL$$

For isotropic distribution of events,  $P(\theta) = \frac{\sin \theta}{2} = \frac{d\Omega}{4\pi} = \frac{2\pi \sin \theta d\theta}{4\pi}$

$$\frac{d\theta}{dL} = \frac{x}{L^2 \sin \theta}$$

Thus  $P(L)_x dL = \frac{x}{2L^2} dL$

Clearly,  $x$  must be  $\leq L$  so there are two cases for the integral over  $x$ , namely,

1.  $L < 1$  then  $P(L) = 2 \int_0^L \frac{x}{2L^2} dx = \frac{1}{2}$   $L < 1$

This 2 takes into account that  $x < L$  near both edges of the slab.

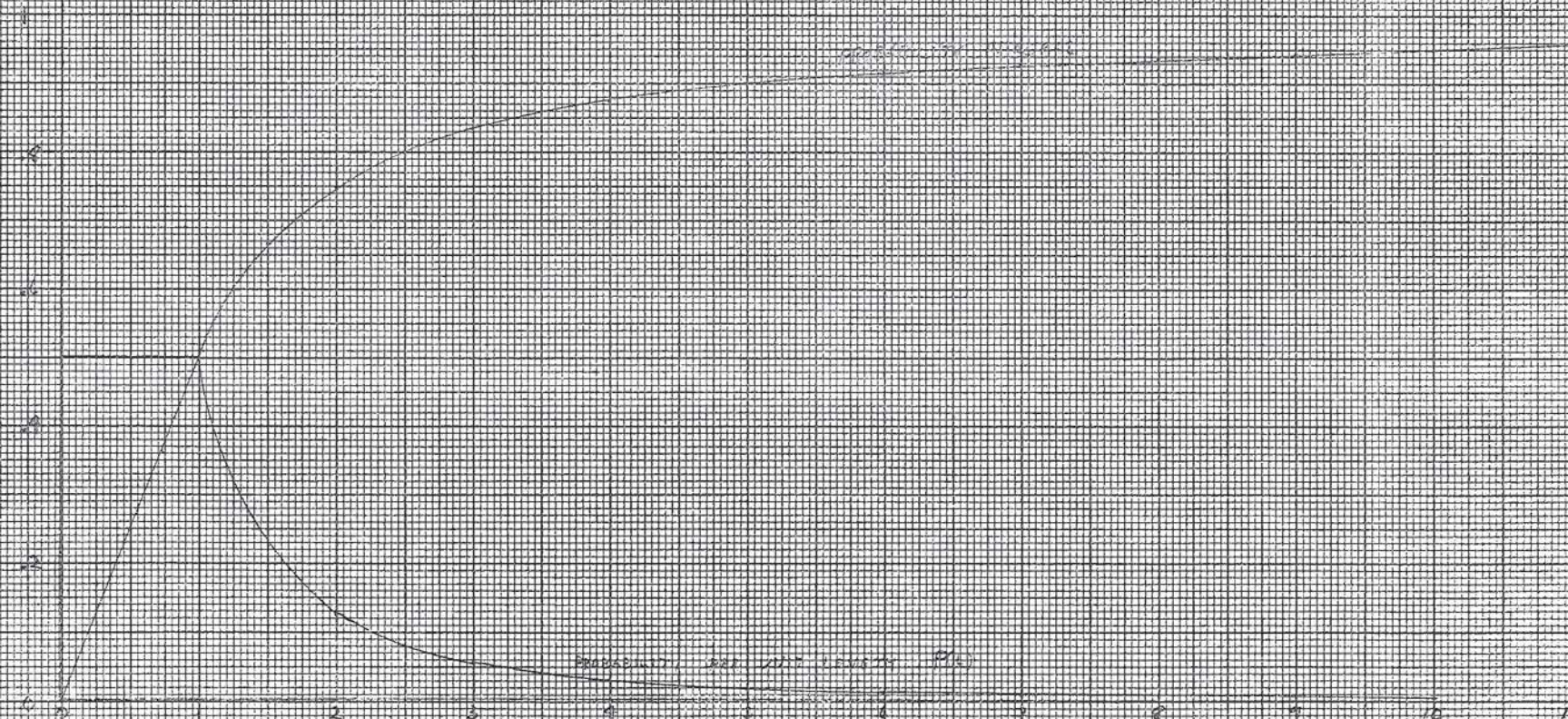
2.  $L \geq 1$   $P(L) = 2 \int_0^1 \frac{x}{2L^2} dx = \frac{1}{2L^2}$   $L \geq 1$

This 2 takes into account lower hemisphere events.



DIRECT ELECTRIC FIELD INTENSITY DISTRIBUTION  
 FOR ISOTROPIC SOURCE IN DIRECTOR OF WAVE  
 LENGTHS AND HEIGHTS GIVEN

HEIGHT (m)



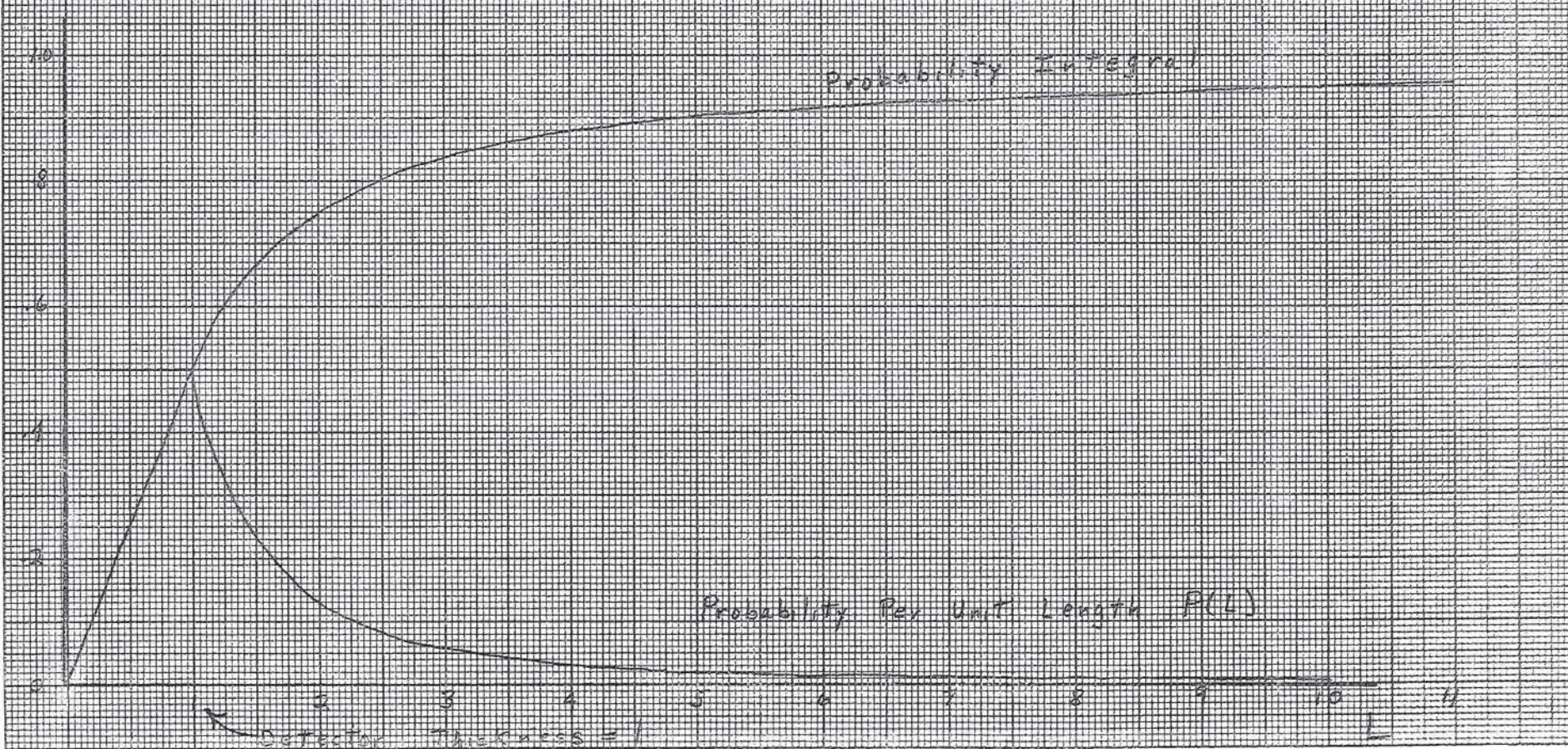
DIRECTOR LENGTH (m)

DIRECTOR HEIGHT (m)



Recall Electron Path Length Distribution  
For Isotropic Events In Detector Of  
Unit Thickness And Infinite Extent

Probability





## Redistribution of Compton Electron Energy loss due to

Escaping from Detector			Energy Loss in Detector											
Initial Recoil Energy	Range $\gamma$ -cm <sup>2</sup>	Detector Stoping power	Energy		Loss in Detector									
			.0	.5	1	1.5	2	2.5	3	3.5	4	4.5	5	
0-.5	0-.243	1.66-1.60	100%	..										
.5-1	.243-.549	1.60-1.47	50	50										
1-1.5	.549-.685	1.47-1.46	50	22	24%									
1.5-2	.685-1.212	1.44-1.47	51	10	7	15%								
2-2.5	1.21-1.55	1.47-1.49	60	20	7	3	10%							
2.5-3	1.55-1.85	1.49-1.50	60	20	7	3	2	8%						
3-3.5	1.85-2.16	1.50-1.52	61	20	6	3	2	1	7%					
3.5-4	2.16-2.47	1.52-1.54	61	20	6	3	2	1	1	6%				
4-4.5	2.47-2.77	1.54-1.55	61	20	6	3	2	1	1	1	5%			

It was assumed that energy loss is constant along the path in the above redistribution.

Energy Interval Mev	Counting Rate per Interval Top Sample	Counting Rate Whole Specimen Mev <sup>-1</sup> Sec <sup>-1</sup>
0-.5	1.959	5.836
.5-1	.102	.408
1-1.5	.0198	.079
1.5-2	.0092	.037
2-2.5	.0035	.014
2.5-3	.0023	.009
3-3.5	.0011	.005
3.5-4	.0010	.004
4-4.5	.0007	.003

Considerations:

Will Escaping Electrons trigger other Detectors, or Scintillator?



Repeat to obtain more data on low end of spectrum. For smaller Recoil Electron Energy

Photon Interval (MeV)	Total Photons counted/sec in interval	Compton Recoil		Electron Energy MeV		Electron Energy MeV		
		.2-.4	.4-.6	.6-.8	.8-1.0	1.0-1.2	1.2-1.4	1.4-1.6
.3-.5	.878	.438	-	-	-	-	-	-
.5-.7	.231	.110	.010	-	-	-	-	-
.7-.9	.101	.028	.040	.005	-	-	-	-
.9-1.1	.068	.014	.015	.024	.002	-	-	-
1.1-1.5	.074	.011	.011	.013	.017	.010	-	-
1.5-2.5	.107	.009	.009	.009	.010	.011	.012	.017
2.5-3.5	.069	.003	.003	.003	.003	.003	.003	.004
3.5-4.5	.095	.002	.002	.002	.002	.002	.002	.002
4.5-5.5	.035	.001	.001	.001	.001	.001	.001	.001

Counting rate for 277 Schwadron for 471 MeV, Sec<sup>-1</sup>

1.605	.616	.091	.057	.035	.027	.018	.029
6.16	.91	.57	.35	.27	.18	.24	

Rate Range Stopping Power Energy lost in detector by escaping electrons

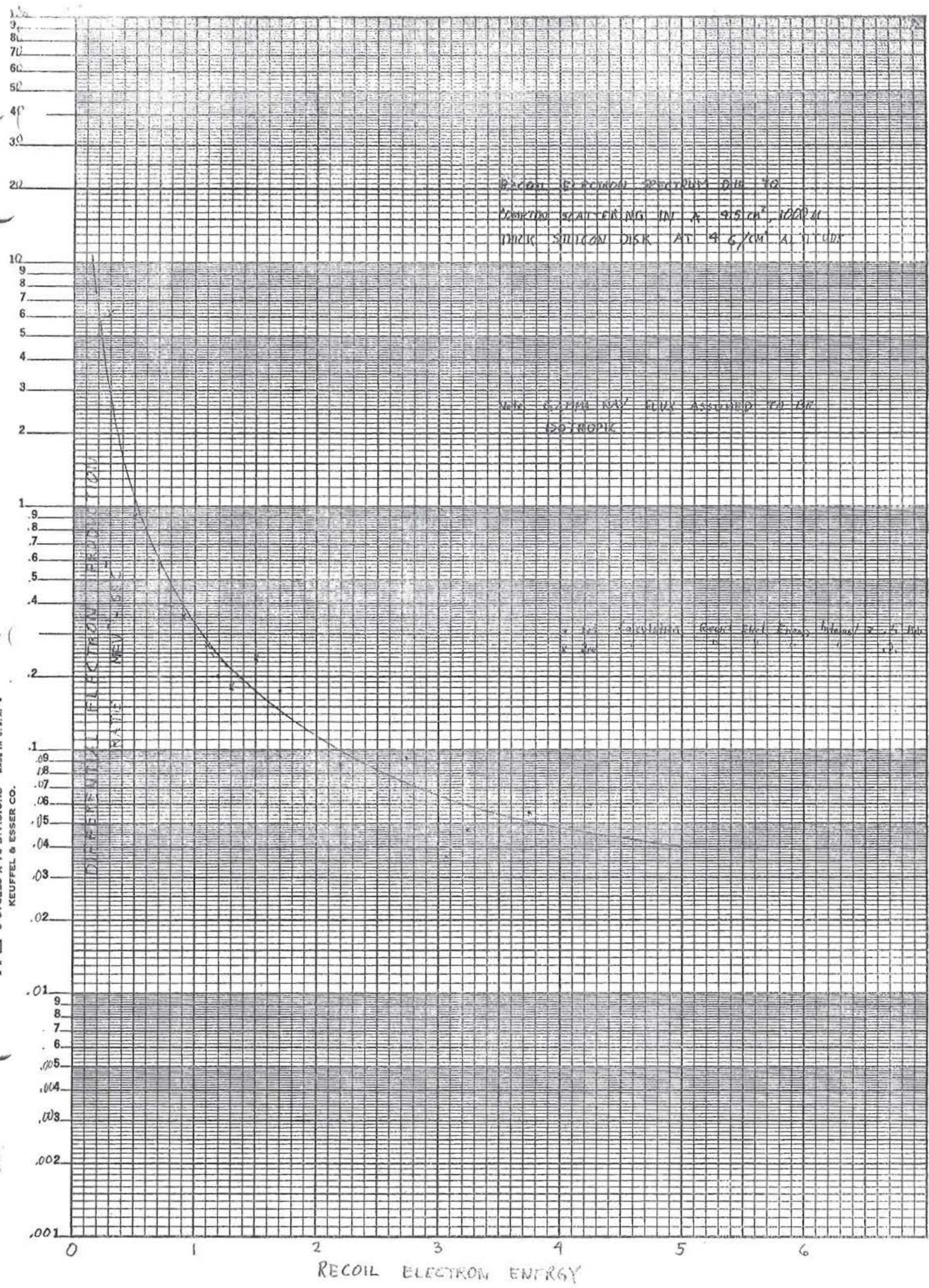
Photon Interval (MeV)	Rate sec <sup>-1</sup>	Range	Stopping Power	Energy lost in detector by escaping electrons										
				<.2	.2-.4	.4-.6	.6-.8	.8-1.0	1.0-1.2	1.2-1.4	1.4-1.6			
.2-.4	.616	.107	1.85	23	77									
.4-.6	.091	.224	1.60	23	24	53								
.6-.8	.057	.351	1.52	25	23	17	35							
.8-1.0	.035	.483	1.48	25	25	16	9	25						
1.0-1.2	.027	.616	1.48	26	25	16	8	5	20					
1.2-1.4	.018	.749	1.46	26	26	15	8	6	3	16				
1.4-1.6	.024	.883	1.46	26	25	16	8	6	3	2	14			

Electron Energy Int. .2-.4 .4-.6 .6-.8 .8-1.0 1.0-1.2 1.2-1.4 1.4-1.6

Counting Rate half spec	.505	.074	.029	.0126	.0067	.0034	.0034
Counting Rate whole spec - MeV <sup>-1</sup> , Sec <sup>-1</sup>	5.05	.74	.29	.126	.067	.034	.034

Contribution of higher incident electron probability, denotes low energy

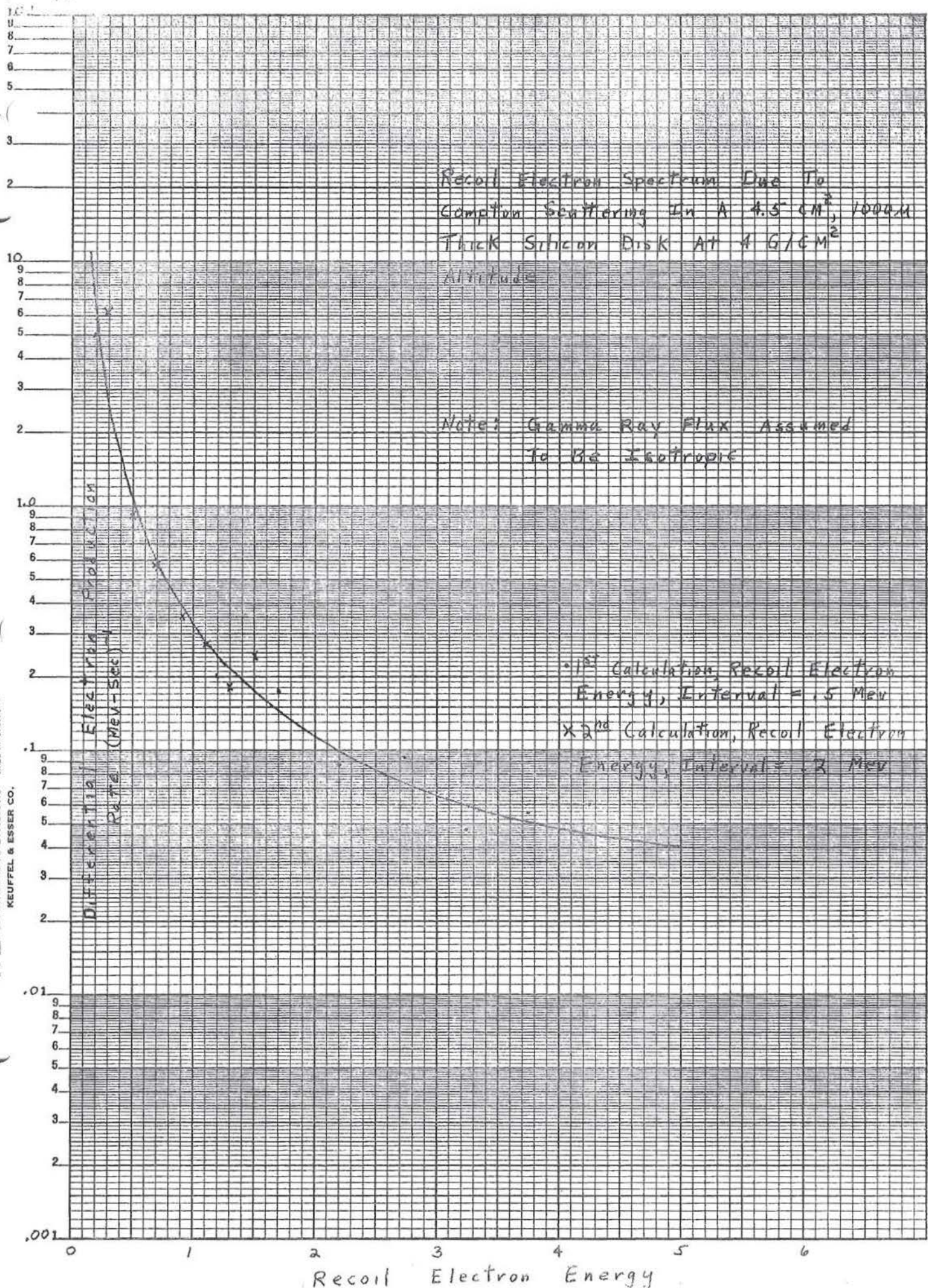




\* see calculation Report and Energy Interval 2.5 to 6.0 keV

RECOIL ELECTRON ENERGY







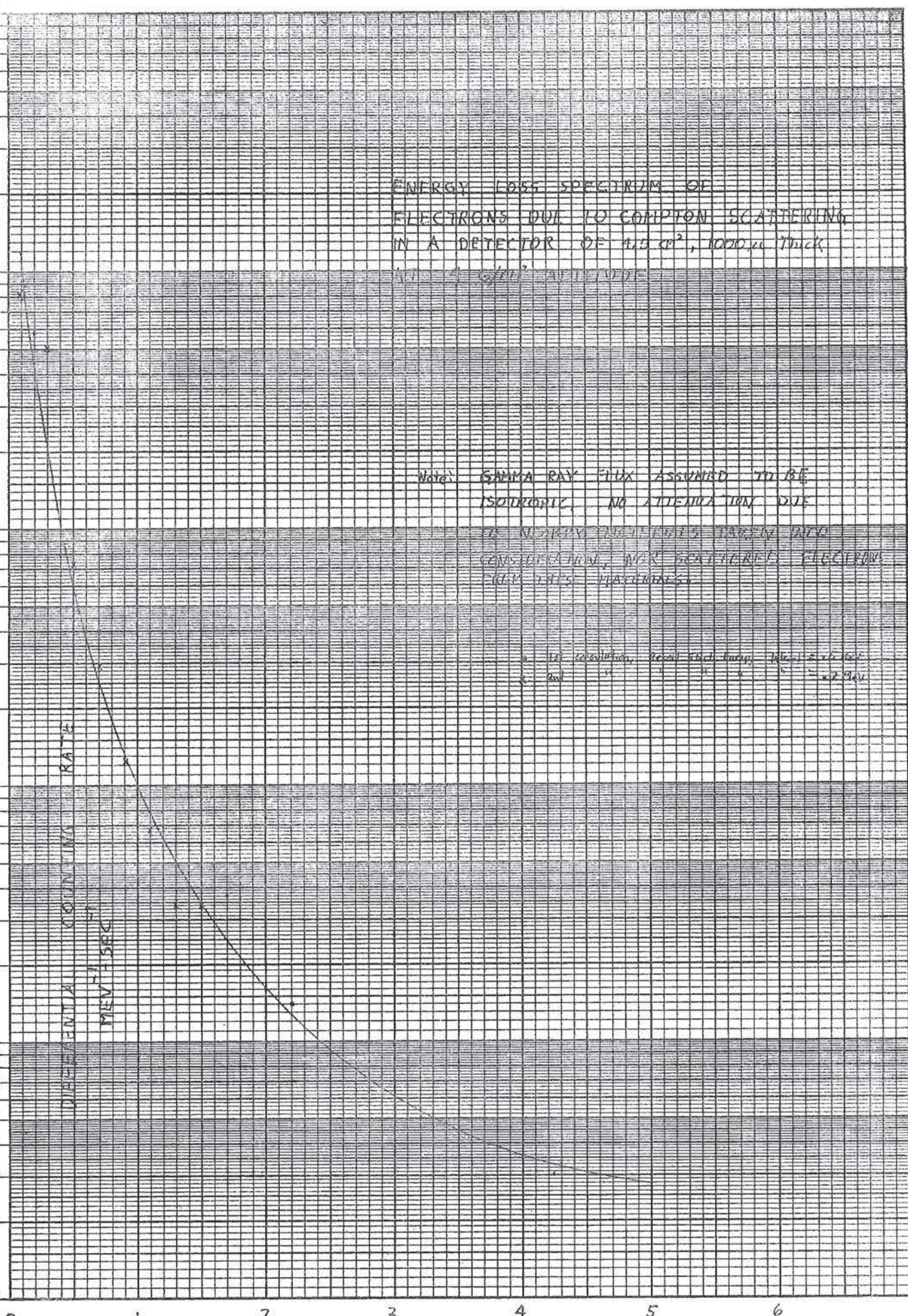
ENERGY LOSS SPECTRUM OF  
ELECTRONS DUE TO COMPTON SCATTERING  
IN A DETECTOR OF 113 CM<sup>2</sup>, 1000 μ thick  
AT 1 G/CM<sup>2</sup> INTENSITY

Note: GAMMA RAY FLUX ASSUMED TO BE  
ISOTROPIC. NO ATTENUATION DUE  
TO SUBSTRATE MATERIALS TAKEN INTO  
CONSIDERATION, NOR SCATTERED ELECTRONS  
FROM THESE MATERIALS.

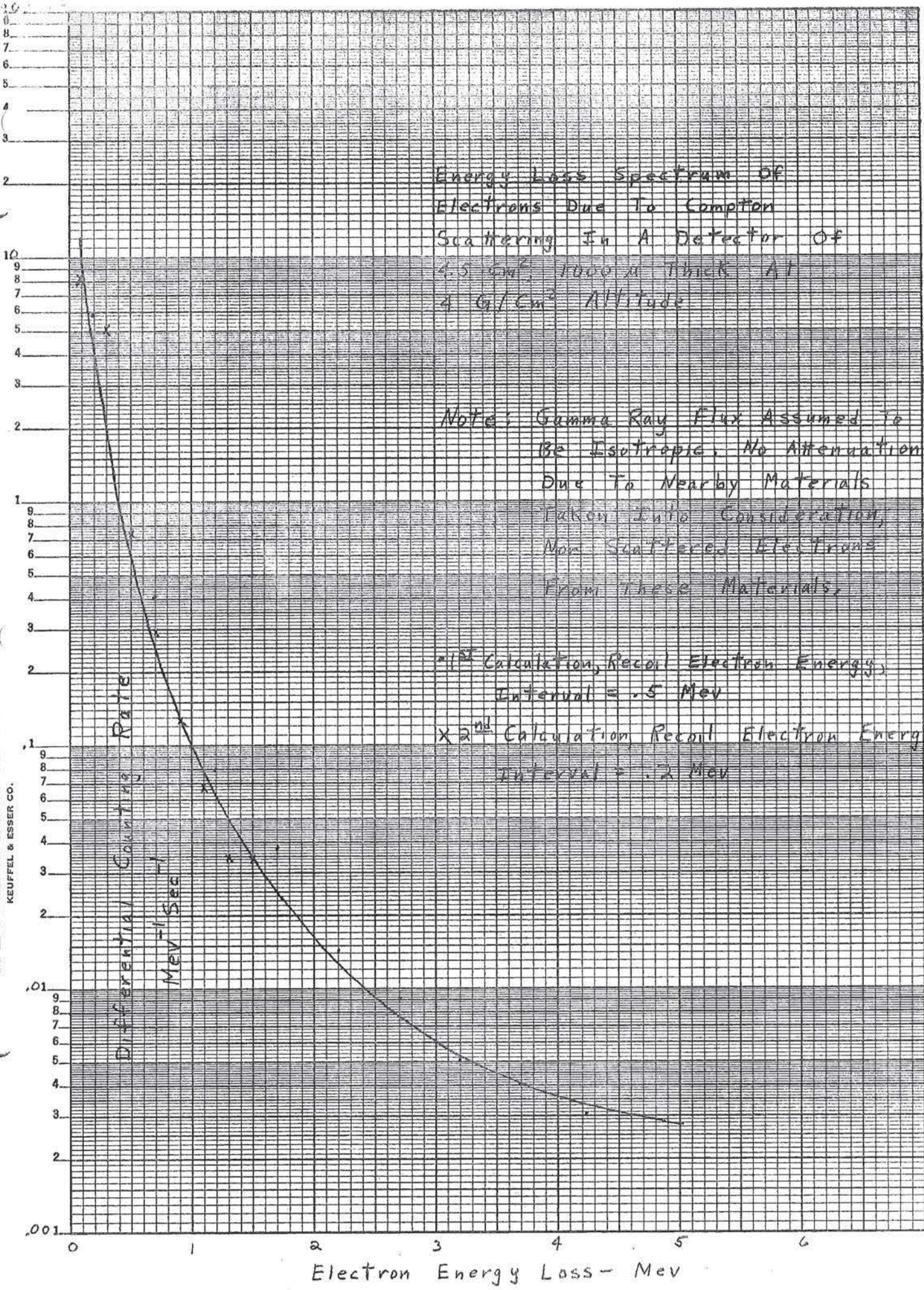
1st resolution, actual peak energy, 1000 μ = 1.02 MeV  
2nd " " " " " " = 1.27 MeV

DIFFERENTIAL COUNTING RATE  
MEV<sup>-1</sup> SEC<sup>-1</sup>

Electron Energy Loss - MeV





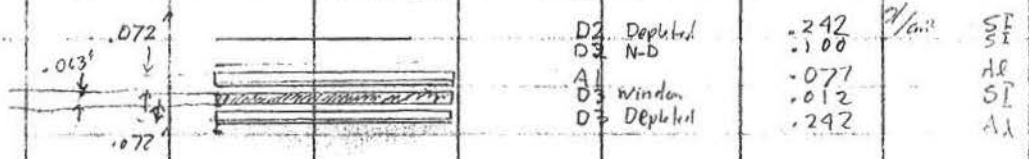




D2 - D3 Events for 25 Compton Electrons

Problem: What are the energy loss spectra in D2 and D3 of electrons scattered by Compton process in U-235?

The Geometry of The detectors is as follows

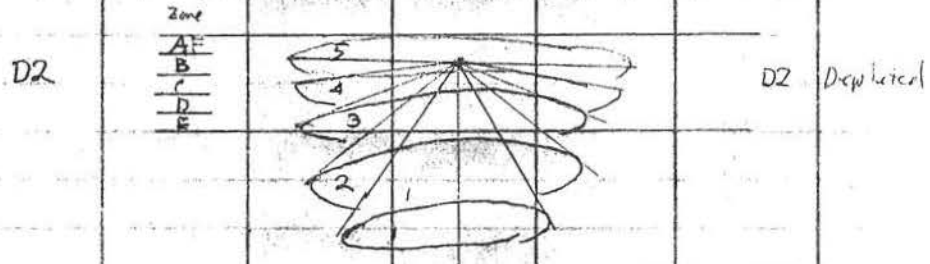


The Compton Recoil Electrons Are scattered isotropically in D2 with known Energy Spectrum. How much energy is lost in D2 & D3 for electrons that lose at least .150 MeV in Both D2 & D3?

Obviously, We are interested in electrons of energy greater than .3 MeV, and owing to the Non-depleted region of D2 and the Absorber Al and Window which Absorb energy, The electrons must at least have .4 MeV to trigger both detectors, but a .4 MeV electron has only .164 cm<sup>2</sup> range which would not get it through the material between detectors.

The Calculation will proceed as follows.

Divide Top detector into 5 Zones of equal depth, and the angular distribution of recoil electrons into 5 zones of equal solid angle (4π/10). Apply the initial recoil spectrum to each of the 20 zones and calculate the energy loss for both detectors. (Ignore 5th Solid Angle Zone)









D2 D3 study (Events for which the energy loss in D2 and D3 is  $\geq 15$  MeV)

Energy loss spectrum D2 for Compton events occurring in D2	Energy loss spectrum D3 for Compton Events occurring in D2	Sum of D2 + D3	Energy	Count MeV <sup>-1</sup>
.15-.3	.15-.3	.0616	.22	.41
.3-.45	.3-.45	.0557	.37	.37
.45-.6	.45-.6	.0508	.51	.39
.6-.75	.6-.75	.0157	.67	.41
.75-.9	.75-.9	.0205	.82	
.9-1.05	.9-1.05	.0027	.90	.077
1.05-1.20	1.05-1.20	.0049		
1.20-1.35	1.20-1.35	.0072	1.20	.039
1.35-1.50	1.35-1.50	.0032	1.42	.021

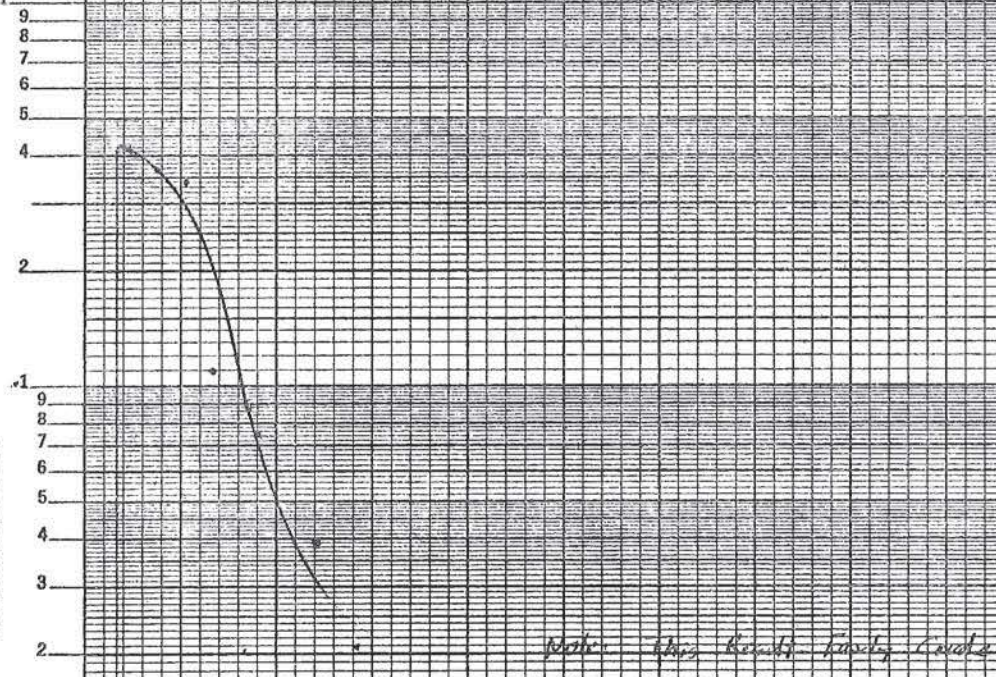
Since Compton scattering may also occur in D3 the scattered electron goes into D2, must add D2 + D3 above to get the total spectrum in either D2 or D3.



KE SEMI-LOGARITHMIC 46 6213  
5 CYCLES X 70 DIVISIONS  
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DIFFERENTIAL COUNTING RATE  
COUNTS PER MINUTE

D2-D3 EVENTS DUE TO  
COMPTON ELECTRONS AT 4.5/CM<sup>2</sup>  
ALTITUDE



Note: this result fairly good

Energy Loss MeV in D2 or D3



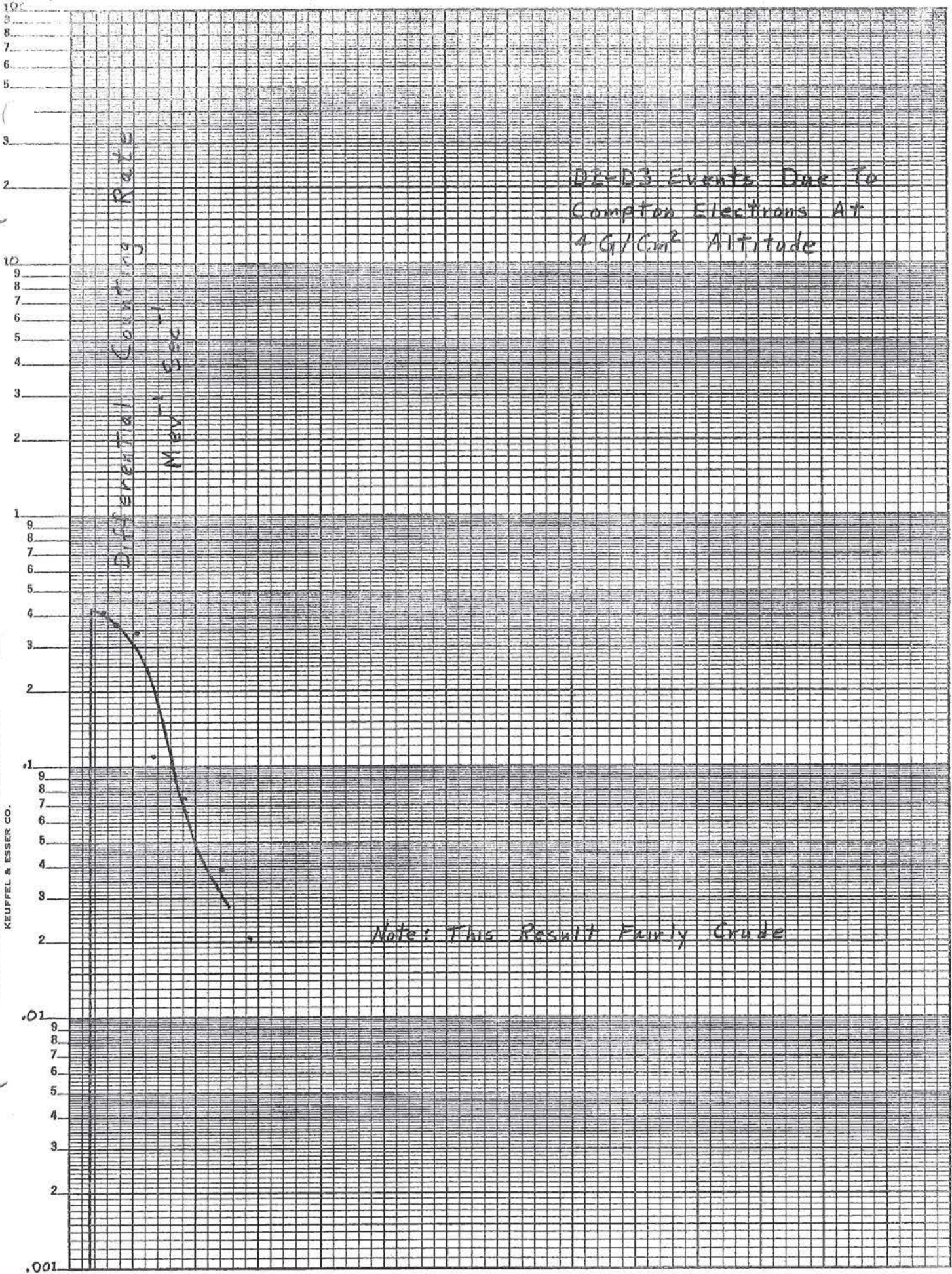
K&E SEMI-LOGARITHMIC 46 6213  
5 CYCLES X 70 DIVISIONS  
MADE IN U.S.A.  
KEUFFEL & ESSER CO.

Differential Counting Rate  
MEV<sup>-1</sup> SEC<sup>-1</sup>

D2-D3 Events Due To  
Compton Electrons At  
4 G/CM<sup>2</sup> Altitude

Note: This Result Fairly Crude

Energy Loss (MEV) In D2 or D3





## References:

### Gamma Ray Spectrum:

Vette, J. I. "Low Energy Gamma Rays Produced in Air and in Lead by Cosmic Rays." JGR 67-5 May, 1961

R. Roshia, J. Labeyrie et al. "Gamma Ray Generation in High Atmosphere", Proc. of Ninth International Conference on Cosmic Rays, Vol. 1.

### Photon Attenuation Data:

NBS Circular 583 and Supplement "X-Ray Attenuation Coefficients 10 KeV to 100 MeV"

### Compton Scattering Data:

NBS Circular 592 "Graphs of the Compton Energy-Angle Relationship and the Klein-Nishina Formula from 10 KeV to 500 MeV"

### Electron Range and Stopping Power Data:

Berger and Seltzer: Tables of Energy Losses and Ranges of Electrons and Positrons, N.S.A. SP 3012, 1964