# Determination of Z, A, and Initial Energy from Voyager HET Events 

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## ABSTRACT

This report describes a method for determining $Z, A$, and the initial energy for a particle traversing a detector with at least two pulse height analyzed $\Delta \mathrm{E}$ measurements. First there will be a brief general discussion of the method, followed by a detailed description of the analysis of Voyager High Energy Telescope (HET) events using this procedure. The three modes covered are HET high gain A-stopping (AS), HET high gain Bstopping (BS), and HET high gain penetrating (PEN). The analysis was done for H and He , including ${ }^{2} \mathrm{H}$ and ${ }^{3} \mathrm{He}$ where possible. A separate report will describe the use of the analysis method on high $Z(>2)$ HET events when that analysis is completed.

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## Introduction

For a particle of a given charge (Z), mass (A), and initial kinetic energy (E), going through a stack of $n$ detectors, a point can be plotted in an $n$-dimensional space, where each coordinate (i) is the energy ( $\Delta \mathrm{E}_{\mathrm{i}}$ ) deposited in the i -th detector. When Z and A are fixed and E is varied, a track is traced out in this n-dimensional space. In reality, events are going to be spread in $n$-space about this track. Ideally, for any event of given $Z$ and A, the most likely initial energy would be that energy corresponding to the point on the theoretical track closest to the point for the event, and the distance ( $r$ ) between the point and the track, with each of the coordinates weighted by their respective fluctuations $\left(\sigma_{\Delta E}\right)$, is a measure of how "good" the event is. Therefore, events on the track can be discriminated from background events by selecting events with $r$ less than a certain value. This value may depend upon energy (i.e. the position along the track) because the fluctuations about the track vary with energy. One should note that certain regions of the track will have to be left out of the analysis due to ambiguities arising from crossing tracks. This crossing of tracks can occur between tracks of different Z and A , and for different energy regions of the same track.

There are also complications due to the shape of the $\mathrm{dE} / \mathrm{dx}$ vs. E curve. After a steep decrease in $\mathrm{dE} / \mathrm{dx}$ with increasing energy, there is a minimum value of the energy loss at the "minimum ionizing energy" and then a slow increase (the relativistic rise). Because of this shape, there is a region of incident energies from slightly less than minimum ionizing on up for which the track is completely ambiguous. This effect is what determines the upper limit for incident energies in this analysis.

In practice, determination of $r$ and $E$ are not this easy. Voyager's solid-state detectors contain "dead layers" in which the energy deposited is not added into the pulse height analysis. This causes an offset in the track for particles which penetrate the dead layer. Because the dead layers are thin $(\sim 150 \mu \mathrm{~m})$, this problem only affects particles in
a small energy range (a few $\mathrm{MeV} / \mathrm{n}$ ), smaller than the final energy resolution. However, these breaks do make the calculation of r and E analytically difficult. Instead, a computationally simpler procedure which roughly approximates this method is used.

If one differentiates between particles which stop in a given detector and particles which fully penetrate the detector, and if dead layers are neglected, then for a given $Z$ and $A$, the energy deposited in the detector is a monotonic function of the initial energy with the $\Delta E$ monotonically increasing with increasing energy for stopping particles, and decreasing with increasing energy for penetrating particles, ignoring particles with incident energies more than the minimum ionizing energy. Therefore, taking each of these cases separately, the value of the initial energy is a direct function of the $\Delta \mathrm{E}$. The Voyager High Energy Telescopes (HETs) have three detectors (or stacks of detectors) which are pulse height analyzed for each event. Figure 1 shows a cross-section of a HET and shows the detectors or stack of detectors which pulse-height analyzed for each of the three modes. This gives three independent values of the initial energy, E, for each event, if Z and A are assumed. Then, the average of the three $\mathrm{E}_{\mathrm{i}}$ (perhaps a weighted average) is approximately the closest E on the theoretical track, and the variance of the $\mathrm{E}_{\mathrm{i}}, \sigma_{\mathrm{E}}^{2}$, as given by $\Sigma\left(\mathrm{E}_{\mathrm{i}}-\mathrm{E}_{\mathrm{svg}}\right)^{2}$ gives a measure of how close the point is to the theoretical track, and so can be used to select events.

One problem with this is the breaks in the $\Delta E$ vs. E curves caused by the dead layers. This was resolved by artificially smoothing the curves so that the $\Delta \mathrm{E}_{\mathrm{i}}$ are always monotonic functions of E . This also allows the analysis to use a previously written fast spline interpolation routine which requires monotonic functions. This smoothing is acceptable because the energy region affected (stopping in the dead layers) is smaller than the final energy resolution.

Another problem arises due to the fact that, in regions where the slope, $\delta \Delta \mathrm{E}_{\mathrm{i}} / \delta \mathrm{E}_{\mathrm{i}}$, is small, a small fluctuation in $\Delta \mathrm{E}_{\mathrm{i}}$ causes a large shift in $\mathrm{E}_{\mathrm{i}}$, and therefore in $\mathrm{E}_{\text {avg }}$ and in the variance, $\sigma_{\mathrm{E}}^{2}$. However, $\mathrm{E}_{\text {avg }}$ can be used to determine three $\Delta \epsilon_{\mathrm{i}}$, the nominal theoretical values for $\Delta \mathrm{E}_{\mathrm{i}}$ given an initial energy $\mathrm{E}_{\text {avg }}$, and another variance, $\sigma_{\Delta E}^{2}$, can be calculated from $\Sigma\left(\Delta \epsilon_{\mathrm{i}}-\Delta \mathrm{E}_{\mathrm{i}}\right)^{2}$ which works well as a measure of how "good" the event is in the region where $\delta \Delta \mathrm{E}_{\mathrm{i}} / \delta \mathrm{E}_{\mathrm{i}}$ is small, but not in the region where it is large. But, by using a combination of both variances, $\sigma_{E}^{2}$ and $\sigma_{\Delta E}^{2}$, the appropriate events can be selected across the entire energy range. The cut using both variances was determined empirically and will be discussed later.

## Selection Of Events And Preliminary Analysis

Several steps are required before calculating Z, A, and E. Figure 2 is a flow chart of these and subsequent steps in this analysis. The Voyager program, vrebox, is used to select the events to be analyzed. There are six different input files, three each for Voyager 1 and Voyager 2: A-stopping, B-stopping, and penetrating. Only high gain events are currently analyzed for H and He , and only events without the caution flag set are included. Also, for Voyager 1, only HET 2 is analyzed due to a PHA problem on HET 1, and on Voyager 2, only HET 1 is used because HET 2 is currently only in low gain mode. Although the other telescopes are usable early in the Voyager flights, the extra geometry is not at this point worth the time required to include their analysis. Wide window (boxes) on the raw pulseheights are also used to discriminate against particles far from the hydrogen and helium tracks. The slant cut is irrelevent, and the guard conditions are discriminated against at a later point in the process for penetrating and are set at $\overline{\mathrm{GI}}$ for A-stopping and B-stopping.

Then the Voyager programs ecal\#_lowz (where \# is the space craft number, 1 or 2 ) are used to convert the three digital PHA values in to $\Delta \mathrm{E}_{\mathrm{i}}$ in MeV using pre-flight calibration data. These programs also add a random number between 0 and 1 to the digital pulse height values in order to smooth out the events in $\Delta \mathrm{E}$ space. This is done because it is easier to visually analyze such factors as event density when the events are spread out. This random number should have a negligible effect on the initial energy value calculated.

The two ecal programs also write a "tel" number for each event, which identifies the particular telescope and the range for each event. All of the information given by the tel number is also in the tag bits, but tel is used to discriminate between events for simplicity.

## Initial Energy Determination

The main step in this analysis involves the three versions of a program "eincal": eincalas for A-stopping events, eincalbs for B-stopping events, and eincalpen for penetrating events. These programs and all the auxilliary files, and the fluxcal programs (which will be described later) are separated into directories according to the mode which they analyze. The three directories are Odin:/usr/erc/voyager/astop, Odin:/usr/erc/voyager/bstop, and Odin:/usr/erc/voyager/pen. The eincal programs take as input the chapter-verse format written out by ecal\#_lowz, and output a chapter-verse format with all of the chapter 41s (individual event data) have been converted into chapter 39 s in which the PHA data has been removed, and values for Z and EIN (initial energy per nucleon) are included for every event. The format of a chapter 39 is as follows:

| Chapter 39 EINCAL Output |  |  |
| :---: | :--- | :--- |
| Byte Number | Length | Name |
| 0 | 2 | KEY = 39 (Chapter Number) |
| 2 | 2 | Spare |
| 4 | 2 | NMVN = Number of Events ( $\leq 40)$ |
| 6 | 2 | Spare |
| 8 | 4 (long) | Itel |
| 12 | 4 (float) | Y1 $\left(\Delta \mathrm{E}_{1}\right) \mathrm{MeV}$ |
| 16 | 4 (float) | $\mathrm{Y} 2\left(\Delta \mathrm{E}_{2}\right) \mathrm{MeV}$ |
| 20 | 4 (float) | $\mathrm{Y} 3\left(\Delta \mathrm{E}_{3}\right) \mathrm{MeV}$ |
| 24 | 4 (float) | $\mathrm{Z}(\mathrm{Charge})$ |
| 28 | 4 (float) | EIN $(\mathrm{MeV} / \mathrm{amu})$ |
| $32 \ldots$ | Repeat last 24 bytes NMVN times (one set per event) |  |

The telescope numbers have also been modified to include such information as guard conditions, particle direction, and mass. The new telescope numbers are:

| tel | HET | Gain | Guard | Direction | Mass |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 300 | 1 | HG | G1* | A-Stop | Lighter Isotope ( ${ }^{1} \mathrm{H}$ or ${ }^{3} \mathrm{He}$ ) |
| 301 | 1 | HG | G1* | A-Stop | Heavier Isotope ( ${ }^{2} \mathrm{H}$ or ${ }^{4} \mathrm{He}$ ) |
| 400 | 1 | HG | G1* | B-Stop | Lighter Issotope ( ${ }^{1} \mathrm{H}$ or ${ }^{3} \mathrm{He}$ ) |
| 401 | 1 | HG | G1* | B-Stop | Heavier Isotope ( ${ }^{2} \mathrm{H}$ or ${ }^{4} \mathrm{He}$ ) |
| 460 | 1 | HG | G1.G2.G3 ${ }^{*}$ | A-Pen | all (isotopes inseparable) |
| 461 | 1 | HG | G1.G2.G3* | B-Pen | all (isotopes inseparable) |
| 470 | 1 | HG | G1.G2.G3 | A-Pen | all (isotopes inseparable) |
| 471 | 1 | HG | G1-G2.G3 | B-Pen | all (isotopes inseparable) |
| 480 | 1 | HG | G1* | A-Pen | all (isotopes inseparable) |
| 481 | 1 | HG | G1* | B-Pen | all (isotopes inseparable) |
| 490 | 1 | HG | G1. G2 ${ }^{\text {a }}$ | A-Pen | all (isotopes inseparable) |
| 491 | 1 | HG* | $\mathrm{G} 1 . \mathrm{G} 2{ }^{*}$ | B-Pen | all (isotopes inseparable) |
| 700 | 2 | HG | G1* | A-Stop | Lighter Isotope ( ${ }^{1} \mathrm{H}$ or ${ }^{3} \mathrm{He}$ ) |
| 701 | 2 | HG | G1* | A-Stop | Heavier Isotope ( ${ }^{2} \mathrm{H}$ or ${ }^{4} \mathrm{He}$ ) |
| 800 | 2 | HG | G1* | B-Stop | Lighter Isotope ( ${ }^{1} \mathrm{H}$ or ${ }^{3} \mathrm{He}$ ) |
| 801 | 2 | HG | G1* | B-Stop | Heavier Isotope ( ${ }^{2} \mathrm{H}$ or ${ }^{4} \mathrm{He}$ ) |
| 860 | 2 | HG | G1.G2.G3 ${ }^{\circ}$ | A-Pen | all (isotopes inseparable) |
| 861 | 2 | HG | G1.G2.G3* | B-Pen | all (isotopes inseparable) |
| 870 | 2 | HG | G1-G2.G3 | A-Pen | all (isotopes inseparable) |
| 871 | 2 | HG | $\mathrm{G} 1 \cdot \mathrm{G} 2 \cdot \mathrm{G} 3$ | B-Pen | all (isotopes inseparable) |
| 880 | 2 | HG | G1* | A-Pen | all (isotopes inseparable) |
| 881 | 2 | HG | G1* | B-Pen | all (isotopes inseparable) |
| 890 | 2 | HG | G1. $\mathrm{G2}^{*}$ | A-Pen | all (isotopes inseparable) |
| 891 | 2 | HG | $\mathrm{G} 1 . \mathrm{G} 2^{*}$ | B-Pen | all (isotopes inseparable) |

These programs ask for an "ascii control file name" which is an ascii file containing a table of points on the theoretical curves for each of the Z and A of interest. The format of each line is: EIN $\Delta E_{1} \Delta E_{2} \Delta E_{3}$. The penetrating files are $v 1 h 2 p$ pein and v 2 h 1 p .ein and contain, in the following order, 27 lines for hydrogen A-penetrating (events coming from the $A$ end of the detector), 27 lines for hydrogen $B$-penetrating, 27 lines for helium4 A-penetrating, and 27 lines helium4 B-penetrating. For A-stopping, the files are v1h2a.ein and v2h1a.ein and contain 27 lines each for ${ }^{1} \mathrm{H},{ }^{2} \mathrm{H},{ }^{3} \mathrm{He}$, and ${ }^{4} \mathrm{He}$. The Bstopping files are $v 1 \mathrm{~h} 2 \mathrm{~b}$.ein and v2h1b.ein and are in the same format as the A-stopping files.

These files were obtained with the program deltae.c which uses Rick Cook's spline fitting to energy loss curves from wrespni.c. This'requires knowing the thicknesses of all layers, including dead layers. Some work has been done on this (see Alan Cumming's Voyager notebooks), and I decided to use those derived thicknesses, which are noticeably different from detector to detector. The thicknesses include a factor of 1.026 which is approximately the average $\sec \theta$ for incident particles.

|  |  | Thickness $(\mu \mathrm{m})$ |  |  |  |
| :---: | :---: | ---: | ---: | ---: | ---: |
| Label | Active | V1HET1 | V1HET2 | V2HET1 | V2HET2 |
|  | no | 50.00 | 50.00 | 50.00 | 50.00 |
| A1 | yes | 150.30 | 150.70 | 146.20 | 151.50 |
| A2 | yes | 149.30 | 149.70 | 142.30 | 150.20 |
| C1 | yes | 3214.00 | 3040.00 | 3035.00 | 3168.00 |
|  | no | 89.30 | 60.60 | 73.90 | 97.50 |
| C2 | yes | 3594.00 | 3435.00 | 3109.00 | 3427.00 |
|  | no | 157.00 | 176.50 | 149.80 | 142.70 |
| C2 | yes | 3280.00 | 3177.00 | 2950.00 | 3427.00 |
| C3 | yes | 3362.00 | 3301.00 | 3018.00 | 3309.00 |
|  | no | 129.30 | 153.00 | 184.70 | 154.00 |
| C3 | yes | 3362.00 | 3302.00 | 3018.00 | 3309.00 |
| C4 | yes | 3000.00 | 3300.00 | 3000.00 | 3000.00 |
|  | no | 155.95 | 153.90 | 152.87 | 176.60 |
| C4 | yes | 3000.00 | 3000.00 | 3000.00 | 3000.00 |
|  | no | 63.60 | 87.21 | 68.74 | 88.80 |
| B2 | yes | 2246.00 | 2200.00 | 2200.00 | 2200.00 |
| B1 | yes | 2238.00 | 2200.00 | 2200.00 | 2200.00 |
|  | no | 87.20 | 60.53 | 88.24 | 76.60 |
|  | no | 64.00 | 64.00 | 64.00 | 64.00 |

To first order, these thickness differences look the same as gain changes in the analog to digital conversion, but these differences affect the ranges of initial energy for which
a particle stops in a detector, and have second order effects in the $\Delta E$ tracks. The second order effects were too small to be seen, so these thicknesses and their appropriate energy ranges were used throughout the procedure. It should be noted, however, that there may be systematic differences between the incoming energy and the calculated value on the order of a few percent.

However, using these thicknesses, and the standard gains from ecal\#_lowz programs, gave a disagreement between the theoretical curve and the center of the distribution of events. This was corrected by using multiplicitive factors in the programs which modify the various $\Delta \mathrm{E}_{\mathrm{i}}$ so that they agree with the theoretical curves. These multiplicitive factors are input into the eincal programs in a gain file given as the required single argument to the programs. These gain files consist of three ascii numbers which are the gains for each of the three $\Delta \mathrm{E}_{\mathrm{i}}$. These gains were obtained empirically by comparing $\Delta \mathrm{E}_{\mathrm{i}}$ vs. $\Delta \mathrm{E}_{\mathrm{j}}$ plots to the theoretical tracks. Figures 3 through 11 show the three $\Delta \mathrm{E}_{\mathrm{i}}$ vs. $\Delta \mathrm{E}_{\mathrm{j}}$ plots for each of the three modes with the theoretical tracks added in. These were generated with the programs eintesta, eintestb, and eintestp which read the ecal tapes and output the first 10000 events into three files, one for each $\Delta E_{i}$ vs. $\Delta E_{j}$ plot. The eintest programs also require a gain file as an argument, so the tracks match up unless the gains have shifted. The gain files are named v\#gainaX.Y, v\#gainbX.Y, and $v$ \#gainpX.Y where \# is the spacecraft number ( 1 or 2 ) and $X$ and $Y$ are respectively the first and last CRS tapes for which these gains are appropriate (e.g. v1gainb1.31 is the V1 B-Stopping gain file for CRS tapes 1 through 31 ). The $\Delta \mathrm{E}_{\mathrm{i}}$ vs. $\Delta \mathrm{E}_{\mathrm{j}}$ plots for all tapes, and their matching theoretical tracks are in my file folders 'HET A-Stopping $H$ and He Final", "HET B-Stopping H and He Final", and "HET Penetrating H and He Final". If the gains are wrong by $\sim 2 \%$ or less, the fluxes are different by only a few percent, and the gains obtained (by eyeball fit) can match the distribution to the theoretical tracks easily by less than $2 \%$.

When a new ecal tape is ready, run the shell script gainbal\#, where \# is the voyager number. The gainbal scripts require one argument, the most recent gain file. The gainbal files are different in the different directories (astop, bstop, and pen), and so this step needs to be done 6 times ( 2 Voyagers $\times 3$ modes). Two plots will be generated by each of these programs. These plots should be labelled with the CRS tape number and the gain file. They then have to be compared to the theorectical tracks by eye. If the theoretical tracks do not match up with the particles distributions, then the gains need to be changed. If the particles are higher (more $\Delta_{E}$ ) than the theoretical tracks then the gains must be decreased, and vice versa. After setting up a new gain file, rerun gainbal to check the new gains.

Yet this was not always enough to get complete agreement. Although these gains successfully match the helium tracks, in some detectors the hydrogen tracks require a slightly different gain. I tried adjusting the thicknesses to see if the second-order effects from incorrect thicknesses could cause this problem, but it seems unlikely that the thicknesses are as far wrong ( $>25 \%$ ) as are necessary to cause these $2 \%$ effects between hydrogen and helium. So ${ }^{+1}$ assume that this is caused by uncorrected non-linearities in the ADCs or shifts in the offsets of the ADCs. Ignoring the possible cause, the correction was empirically made in the input files to the eincal programs. Two of the theoretical $\Delta \mathrm{E}_{\mathrm{i}}$ had all of their hydrogen values decreased by hand by $2 \%$ to bring them into line with the event distributions: B1 in the file v1h2b.ein (B-Stopping), and B1 in the file v2h1b.ein (B-Stopping).

It should be noted that fragmentation taking place in the first detector is visible in the $\Delta E_{3}$ vs. $\Delta E_{2}$ plots. For Example, in Figure 4 you can clearly see the secondary deuterium and tritium which is not visible in Figure 3. These events will be spead out in the $\Delta \mathrm{E}_{1}$ coordinate and so will be off the theoretical track in the three dimensional space used in this analysis.

Since the exact analysis varied between the three different event modes, due to both program evolution and the tailoring of the programs to each of the modes, they will be treated separately in the following discussion of event discrimination and initial energy determination,

## A Stopping

The program eincalas analyzes $\mathrm{H},{ }^{2} \mathrm{H},{ }^{3} \mathrm{He}$, and ${ }^{4} \mathrm{He}$ for A-stopping events which penetrate as far as detector Cl (three-parameter events). In the analysis, $\Delta \mathrm{E}_{1}=\Delta \mathrm{E}_{\mathrm{A}}$ (in MeV ), $\Delta \mathrm{E}_{2}=\Delta \mathrm{E}_{\mathrm{A} 2}$ (in MeV ), and $\Delta \mathrm{E}_{3}=\Delta \mathrm{E}_{\mathrm{C} 1+\mathrm{C} 2+\mathrm{C} 3}$ (in MeV ). It first selects events by telescope (tel) number such that they are A-stopping: $300 \leq$ tel $<400$ (HET 1) or $700 \leq$ tel $<800$ (HET 2), and it also eliminates noise in low bins by requiring $\ln \left(\Delta \mathrm{E}_{1}\right)>\mathbf{- 2 . 3}$, $\ln \left(\Delta \mathrm{E}_{2}\right)>-2.3$, and $\ln \left(\Delta \mathrm{E}_{3}\right)>0$. Two-parameter events are eliminated by requiring that $C 1$ triggers: $((\operatorname{tag} \& 04000) \neq 0)$. Since the hydrogen and helium tracks don't cross, the program does completely separate analyses of the two regions around the hydrogen and helium. Events with $0.2<0.58 \times \ln \left(\Delta \mathrm{E}_{3}\right)+\ln \left(\Delta \mathrm{E}_{2}\right)<2.4$ are possible hydrogen, and events with $0.58 \times \ln \left(\Delta \mathrm{E}_{3}\right)+\ln \left(\Delta \mathrm{E}_{2}\right) \geq 2.4$ are possible helium. If $0.58 \times \ln \left(\Delta \mathrm{E}_{3}\right)+\ln \left(\Delta \mathrm{E}_{2}\right)<0.2$ then the events are background.

For the hydrogen section, three initial energies, $\mathrm{E}_{\mathrm{i}}$, are calculated from the table of $E_{i}$ vs. $\Delta E_{i}$ read in from the input file Actually the table is stored in the program as $\ln \left(\mathrm{E}_{\mathrm{i}}\right)$ vs. $\ln \left(\Delta \mathrm{E}_{\mathrm{i}}\right)$ ) for the use of the cubic spline routine which does the interpolation.

Then the average initial energy, $\mathrm{E}_{\mathbf{2 v g}}$ and the variance, $\sigma_{\mathrm{E}}^{2}$, are calculated from:

$$
E_{2 v g}=\sum_{i=1}^{3} \frac{E_{i}}{3} \quad \sigma_{E}^{2}=1000 \times \sum_{i=1}^{3} \frac{\left(E_{2 v g}-E_{i}\right)^{2}}{E_{\mathrm{vvg}}^{2}}
$$

$\mathrm{E}_{\mathrm{svg}}$ is then used to calculate the $\Delta \epsilon_{\mathrm{i}}$ from the inverse tables, $\Delta \mathrm{E}_{\mathrm{i}}$ vs. $\mathrm{E}_{\mathrm{i}}$, and the variance $\sigma_{\Delta E}^{2}$ is obtained from:

$$
\sigma_{\Delta E}^{2}=1000 \times \sum_{\mathrm{i}=1}^{3} \frac{\left(\ln \left(\Delta \mathrm{E}_{\mathrm{i}}\right)-\ln \left(\Delta \epsilon_{\mathrm{i}}\right)\right)^{2}}{\mathrm{E}_{\mathrm{zvg}}^{2}}
$$

The program uses $\ln \left(\Delta \mathrm{E}_{\mathrm{i}}\right)$ and $\ln \left(\Delta \epsilon_{\mathrm{i}}\right)$ because that is what is returned by the interpolation subroutine. Figure 12 shows a plot of $\sigma_{E}^{2}$ vs. $\sigma_{\Delta E}^{2}$ in which it can be seen that the edge of the helium events is diagonal on the log-log plot. From this came the equation used for $r$ which approximates the distance to the theoretical track: $r=0.6 \times \ln \left(\sigma_{\Delta E}^{2}\right)+\ln \left(\sigma_{E}^{2}\right)$. Figure 13 shows the histograms of $r$ for ${ }^{4} \mathrm{He}$. For both hydro gen and helium the distance, $r$, is calculated independently for both isotopes. To be considered a good event, $r$ must be less than the $r$ obtained for the other isotope, and less than an empirically derived value, $\mathrm{r}_{\mathrm{max}}$, which is 4.0 for the hydrogen isotopes and 0.5 for the helium isotopes. Although Figure 13 appears to show that the cut of $\mathrm{r}<0.5$ is discriminating against good ${ }^{4} \mathrm{He}$ events, Figure 12 shows that this value is a reasonable cut, and indeed, to accept any events with larger $r$ would prevent ${ }^{3} \mathrm{He}$ from being resolved. Good events are output with their correct $Z$ value, and other events are assigned $Z=0$. In Figures 14 through 19, the events plotted in Figures 3, 4, and 5 are separated into good events (Z, A, and EIN determined) and background.

## B Stopping

For B-stopping events, the program eincalbs looks almost identical to eincalas. The values of cuts are changed, and in this case $\Delta \mathrm{E}_{1}=\Delta \mathrm{E}_{\mathrm{B} 1}, \Delta \mathrm{E}_{2}=\Delta \mathrm{E}_{\mathrm{B} 2}$, and $\Delta \mathrm{E}_{3}=\Delta \mathrm{E}_{\mathrm{C} 2+\mathrm{C} 3+\mathrm{C}_{4}}$. If $4.6 \geq \ln \left(\Delta \mathrm{E}_{3}\right)+1.8 \times \ln \left(\Delta \mathrm{E}_{2}\right)$, or $\Delta \mathrm{E}_{\mathrm{i}}<1 \mathrm{MeV}$ for any of the three detectors, the events are considered background. If $4.6<\ln \left(\Delta \mathrm{E}_{3}\right)+1.8 \times \ln \left(\Delta \mathrm{E}_{2}\right)<7.5$, the events are nominally hydrogen, otherwise they are nominally helium.

Again $r$ is calculated by $r=0.6 \times \ln \left(\sigma_{\Delta E}^{2}\right)+\ln \left(\sigma_{E}^{2}\right)$ and the $r_{\text {max }}$ is -5.0 for hydrogen and -8.0 for helium. Background events have been assigned $Z=0$ as in A-Stopping. Figures 20 through 25 show the events of Figures 6, 7, and 8 separated into good events and background.

## Penetrating

For several reasons the penetrating mode is much more difficult to analyze than either A-stopping or B-stopping. One problem is the bidirectional nature of this mode. At the lower energies, the slowing down of the particle is noticeable and so the direction of the event can be determined, but at higher energies the width of the track combines events from both the A side (particles which hit detector A1 first) and the B side (particles which hit B1 first). Also, for the penetrating mode the isotopes of both hydrogen and helium are indistinguishable, because the separation of the tracks is smaller than the width. Add to this the fact that the helium high-energy (minimum ionizing) particles blot out a noticeable portion of the hydrogen spectrum, and the penetrating procedure becomes much more difficult and complicated than either of the stopping modes. However, of the three, this is the one which was attempted first, and this is evident in the software. Several improvements could be made to the program eincalpen, and there are large sections of code which could be left out because they were not successful enough to allow the particles involved to be used. For example, calculations are made for particles which stop in the third detector, but this never worked very well, and the energy range (a few MeV ) was small compared to the hundreds of MeV over which the analysis is sucessful for fully penetrating particles.

Penetrating high-gain events have tel numbers of 450 (HET 1) or 850 (HET 2), as given by the ecal program. The first thing that the eincalpen program does is to modify tel to include information about the guard conditions: tel $=460$ ( 860 ) is G1. $\mathrm{G} 2 \cdot \overline{\mathrm{G} 3}$, tel $=470(870)$ is $\mathrm{G} 1 \cdot \mathrm{G} 2 \cdot \mathrm{G} 3$, tel $=480(880)$ is $\overline{\mathrm{G} 1}$, and tel $=490(890)$ is $\mathrm{G1} \cdot \overline{\mathrm{G} 2}$. Then events with tels of $480,490,880$, or 890 and with ( $0.9<\Delta \mathrm{E}_{1}$ and $0.9<\Delta \mathrm{E}_{2}$ ) and $\left(\Delta E_{1}<18\right.$ or $\left.\Delta E_{2}<24\right)$ are accepted and split into four regions of analysis, and several regions in which analysis is impossible (such as where the H and He tracks are mixed).

The region with (e1 > 6.0) and (e3 < 60.0) and ( $(9.091 \times \mathrm{el}-\mathrm{e} 3)>18.182$ ) contains the low energy ( $81-89 \mathrm{MeV}$ ) penetrating hydrogen coming from the A direction. If $\sigma_{\Delta E}^{2}<0.02$ and $\sigma_{\mathrm{E}}^{2}<0.125$ then the particles are accepted.

The $B$ direction low energy hydrogen region is (e2>8.0) and (e3 $<60.0$ ) and $((6.071 \times \mathrm{e} 2-\mathrm{e} 3)>12.143)$. Once again the variance cuts are $\sigma_{\Delta E}^{2}<0.02$ and $\sigma_{\mathrm{E}}^{2}<0.125$.

The region in which $1.211 \times \ln \left(\Delta E_{1}\right)+\ln \left(\Delta E_{2}\right)>3.44$ contains helium events with initial energy between about $80 \mathrm{MeV} / \mathrm{n}$ and $650 \mathrm{MeV} / \mathrm{n}$. The EIN for the direction which gives the lowest value of the variance, $\sigma_{\mathrm{AE}}^{2}$, is assumed to be correct. This is acceptable because in the region where the A and B directions are indistiguishable the difference in EINs calculated for the two directions is smaller than the energy resolution. For the penetrating program, the two variances, $\sigma_{\mathrm{E}}^{2}$ and $\sigma_{\mathrm{AE}}^{2}$, are given by:

$$
\sigma_{E}^{2}=\sum_{i=1}^{3} \frac{\left(\mathrm{E}_{2 v g}-\mathrm{E}_{\mathrm{i}}\right)^{2}}{\mathrm{E}_{2 \mathrm{vg}}^{2}} \quad \sigma_{\Delta E}^{2}=\sum_{\mathrm{i}=1}^{3} \frac{\left(\Delta \mathrm{E}_{\mathrm{i}}-\Delta \epsilon_{\mathrm{i}}\right)^{2}}{\mathrm{E}_{2 v \mathrm{~V}} \times \Delta \mathrm{E}_{\mathrm{i}}}
$$

which are slightly different from the stopping programs. Then, if $75 \mathrm{MeV} / \mathrm{n}<\mathrm{E}_{\mathrm{avg}}<$ $650 \mathrm{MeV} / \mathrm{n}$, and $\sigma_{\mathrm{E}}^{2}<1$ and $\sigma_{\Delta E}^{2}<0.005$, the events are taken to be good helium.

The boundaries for high energy hydrogen events (called high energy because their intitial energies are above the energies for which minimum-ionizing helium mixes with the hydrogen track) are $0.759<\ln \left(\Delta \mathrm{E}_{1}\right)+1.012 \times \ln \left(\Delta \mathrm{E}_{2}\right)<2.602$ and $\Delta \mathrm{E}_{3}<23$. Although the variance $\sigma_{\mathrm{E}}^{2}$ is calculated the same as for helium, $\sigma_{\Delta E}^{2}$ is now calculated as:

$$
\sigma_{\Delta E}^{2}=1000 \times \sum_{i=1}^{3} \frac{\left(\Delta E_{i}-\Delta \epsilon_{\mathrm{i}}\right)^{2}}{\mathrm{E}^{2}}
$$

with E being the minimum of $\mathrm{E}_{\text {avg }}$ or 375 . This was done to prevent events having ridiculously high $\mathrm{E}_{\mathrm{avg}}$ from having small $\sigma_{\Delta \mathrm{E}}^{2}$. Acceptable high energy hydrogen events have $110 \mathrm{MeV}<\mathrm{EIN}<375 \mathrm{MeV}$ and $\sigma_{\mathrm{E}}^{2}<1$ and $\sigma_{\Delta E}^{2}<0.30$. It was later determined that the energy range $110 \mathrm{MeV}<E I N<150 \mathrm{MeV}$ suffers from too much helium contamination, and so these particles are left out of the flux calculations. Figures 26 to 31 show the particles selected and not selected from Figures 9, 10, and 11. Note that although the directional information is included in the tel number for every event (see table on page 5), it is only meaningful at the lower energies of the analysis, and only used in the 81 89 MeV hydrogen energy interval.

## Flux Calculations

The three programs which calculate and plot the fluxes, fluxcala (AS), fluxcalb (BS), and fluxcalp (PEN), are mostly identical. They are run with the shell files runfluxa, runfluxb, and runfluxp which correctly pass the output to the plotter. The only important difference is which of the rates are used to calculated the flux. Since not all events of a given mode are pulse height analyzed, the number of events in a given energy bin has to be divided by a calculated livetime, not just the actual time. The formula for flux, F , (in particles $/\left(\mathrm{m}^{2}\right.$ sec sr $\left.\mathrm{MeV} / \mathrm{nuc}\right)$ is given by:

$$
F=\frac{N \cdot \text { Rate }}{\mathrm{A} \Omega \cdot \Delta E \cdot t \cdot P H A}
$$

where $N$ is the number of analyzed events in the energy bin ( $\mathrm{E}_{\text {min }}<\mathrm{E}<\mathrm{E}_{\text {max }}$ ), Rate is the total rate count for that mode as determined by a dead time free rate scalar, $\mathrm{A} \Omega$ is the geometry factor ( $\mathrm{m}^{2} \mathrm{sr}$ ), $\Delta \mathrm{E}$ is $\mathrm{E}_{\mathrm{max}}-\mathrm{E}_{\mathrm{min}}{ }^{-}(\mathrm{MeV} / \mathrm{nuc})$, t is the total time of detector operation (sec), and PHA is the PHA count for that mode. The major difference between
the three programs is that each uses different values (locations) for Rate, $t$, and PHA. These values come from the chapter 48 rate summaries. If you assign the pointer value returned by doing a getvis(1) on chapter 48 to a long (integer) pointer the offsets of the rate data are:

| Mode <br> Rate Label | A-Stopping <br> AS |  | B-Stopping <br> BSp |  | Penetrating <br> PENH |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | HET 1 | HET 2 | HET 1 | HET 2 | HET 2 | HET 2 |
| Rate | 41 | 83 | 47 | 89 | 51 | 93 |
| t | 42 | 84 | 48 | 90 | 52 | 94 |
| PHA | 361 | 368 | 364 | 371 | 367 | 374 |

These three programs plot each energy bin on a log-log plot of flux vs. energy/nuc, complete with statistical error bars in flux calculated from the number of particles in the bin, and error bars in energy which give the width of the bin.

The fluxcal programs read an input file with the following format:

| NZ | The number of plots (one for each isotope) |
| :---: | :---: |
| Plot Title | \| These twelve lines are repeated for each isotope |
| Z | 1 Charge |
| NBIN | \| Number of energy bins |
| $\mathrm{Z}_{\text {min }} \mathrm{Z}_{\text {max }}$ | \| Charge range |
| $\mathrm{Hiflag}_{1} \cdots \mathrm{Hiflag}_{\text {Nbin }}$ | \| Flags for HET 1 ( $0=$ don't use, $1=$ use ) |
| $\mathrm{H} 2 \mathrm{flag}_{1} \cdots \mathrm{H}^{\text {frlag }}$ NBEN | \|Flags for HET $2(0=$ don't use, $1=$ use $)$ |
| $\Omega 1_{1} \cdots \cdots 1_{\text {NBIN }}$ | \| Geometry factors, HET $1\left(\times 10^{5}\right)$ |
| $\Omega 2_{1} \cdots \cdots 2_{\text {NBIN }}$ | \| Geometry factors, HET $2\left(\times 10^{5}\right)$ |
| $\operatorname{Emin}_{1} \cdots \operatorname{Emin}_{\text {NBIN }}$ | \| Minimum energy for energy bins |
| $\operatorname{Emax}_{1} \cdots \operatorname{Emax}_{\text {NBIN }}$ | \| Maximum energy for energy bins |
| NTEL | \| Number of acceptable telescope values |
| $\mathrm{Tel}_{1} \cdots \mathrm{Tel}_{\text {NTEL }}$ | \| Acceptable telescope values |
| ... | Repeat NZ times |

The geometry factors used are calculated for particles which stop in the center of the range of stopping locations for that energy bin. In calculating these geometry factors, the areas of the detectors are multiplied by a factor, 0.94 , to allow for the fact that $6 \%$ of the area of the detectors, at the edge, gives an incorrect value of of $\Delta \mathrm{E}$ due to edge
effects.
These programs also ask for the number of time periods to be used and the beginning and end of these time period given as volume numbers. This allows the removal of time periods in which rate spikes appear, which would invalidate the calculation of the fluxes. Figure 32 shows what the output of these programs looks like, in this case fluxcalp (PEN) for Voyager 1 helium. Figure 33 is the associated dump of all the values. In order to make it easier to get plots with all modes included (AS, BS, and PEN), the programs also output files which contain five values for each energy bin: the flux, and the endpoints of the error bars. One of these files is generated for each $Z$, and these files can be combined to give an overall plot with the program allflux. The files have names such as v1has.h.87a, v2hbs.he4.86c, and v2hpn.he3.85g depending upon Voyager, mode, Z, A, and time period.

There is a program allflux.c which is run by the shell file, runallflux, which combines the output files from the fluxcal programs onto one plot. It can combine up to 11 files, using a different plot character for each file. It first asks for a spacecraft number ( 0 means both) and a title for the plot. Then it requests the number of time periods of interest, and beginning and end volume numbers for each interval. It never actually uses these values, it just writes them to the plot so this information is not lost. Then it asks for the number of files to be combined and the names of each file. Figures 34 and 35 show the output for this program.

## Conclusion

Although the programs work well as they stand now, there are still quite a few improvements which could be made. The program eincalpen was done first and so lacks some of the refinements of the later programs. Its selection cuts should be made diagonal in $\sigma_{E}^{2}$ and $\sigma_{\Delta E}^{2}$ like the A-stopping and B-stopping programs. Also, because of the large number of events in penetrating mode, the energy bins used in the fluxcalp program have been made quite small, smaller, in fact, than the actual energy resolution possible from the width of the the distribution, so that the energy resolution which appears on the plots is misleading.

All three programs still suffer from some common faults. The normalization of the variances is slightly different from mode to mode, and it is not exactly correct in any of them. They should be normalized by the width of their distribution, which is a function of energy. The normalization used only approximates these widths. Also, the calculation of $E_{\text {avg }}$ should probably be a weighted mean (weighted by the slope of the EIN vs. $\Delta E$ curve, and the width of the distribution). The detector thicknesses and PHA gains are
also still in question.
Yet all of these problems are at the couple of percent level or less and the programs still appear to work quite well, and so these problems have been left for future work. For more information on this work, there is a file folder titled "Determination of $\mathrm{Z}, \mathrm{A}$, and Initial Energy from Voyager HET Events" which contains copies of all programs, and their input and outputt files. There are also many other file folders filled with intermediate steps and test output. These can usually be referenced by the descriptions in my Voyager lab notebooks volumes I and II.

- Figure 1 cross-section of HET


## Pha headout:



|  |  |  |
| :---: | :---: | :---: |
|  |  |  |
| AS |  |  |
| PHA3 | BS | PHA3 |
| PHAZ PHA3 | PHAZ | PHAZ |
| PHA1 | PHAI | PHAI |



Figure 2
$Z, A_{s}$ and Initial Energy Determination

allying combines flux output files
runallfluy sends ont/at to plotter (Imogen)

Figure 3


Figure 4


Figure 5


Figure 6


Figure 7


Figure 8


Figure 9


Figure 10


Wed Aug 17 13:02:18 1988










Tue Aug 16 21:13:18 1988





Tue Aug 16 21:17:52 198日 File: eincalbs.outs









## Figure 33

Mon Nov 2 21:53:33 1987


| Livetime Parameters | $:$ | 1 | 2 |
| :---: | :---: | :---: | :---: |
| Rate Counts | : | 173530 | 167814 |
| Times | : | 850326 | 849960 |
| PHA Counts | : | 77699 | 75818 |
| Livetimes | : | 3.81e+05 | $3.84 e+05$ |
| Geometry factors | : | 1.10e-04 | 1.10e-04 |
|  |  | 1.10e-04 | $1.10 e-04$ |
|  |  | 1.09e-04 | 1.09e-04 |
|  |  | 1.08e-04 | 1.08e-04 |
|  |  | 1.00e-04 | 1.00e-04 |
|  |  | $9.28 \mathrm{e}-05$ | $9.28 \mathrm{e}-05$ |
|  |  | $7.86 \mathrm{e}-05$ | 7.86e-05 |

Figure 34


```
Overall time period : 1987: 157.00-1987: 209.00
Intervals : 365569 370560 1987: 257.00-1987: 209.00
```

Vihas.h.87d -

| Energy | Flux | Lower | Upper |
| ---: | :---: | :---: | :---: | :---: |
| $6.7-9.7$ | $1.27 e-01$ | $9.54 e-02$ | $1.67 e-01$ |
| $9.7-12.7$ | $1.74 e-01$ | $1.37 e-01$ | $2.20 e-01$ |
| $12.7-18.7$. | $1.79 e-01$ | $1.52 e-01$ | $2.10 e-01$ |
| $18.7-24.7$ | $2.89 e-01$ | $2.55 e-01$ | $3.27 e-01$ |
| $24.7-35.0$ | $4.44 e-01$ | $4.11 e-01$ | $4.80 e-01$ |
| $35.0-45.4$ | $4.40 e-01$ | $4.05 e-01$ | $4.77 e-01$ |
| $45.4-60.4$ | $5.41 e-01$ | $5.07 e-01$ | $5.78 e-01$ |

vihbs.h.87d

| Energy | Flux | Lower | Upper |
| :---: | :---: | :---: | :---: |
| $29.1-37.9$ | $3.76 e-01$ | $3.50 e-01$ | $4.03 e-01$ |
| $37.9-46.7$ | $5.21 e-01$ | $4.89 e-01$ | $5.54 e-01$ |
| $46.7-61.4$ | $5.92 e-01$ | $5.63 e-01$ | $6.22 e-01$ |
| $61.4-74.0$ | $7.69 e-01$ | $7.30 e-01$ | $8.10 e-01$ |

vihpa.h.87d

| Energy | Flux | Lower | Upper |
| :---: | :---: | :---: | :---: |
| $150.0-175.0$ | $1.69 \mathrm{e}+00$ | $1.63 \mathrm{e}+00$ | $1.74 \mathrm{e}+00$ |
| $175.0-200.0$ | $1.84 \mathrm{e}+00$ | $1.79 \mathrm{e}+00$ | $1.90 \mathrm{e}+00$ |
| $200.0-225.0$ | $1.92 \mathrm{e}+00$ | $1.86 \mathrm{e}+00$ | $1.98 \mathrm{e}+00$ |
| $225.0-250.0$ | $1.93 \mathrm{e}+00$ | $1.87 \mathrm{e}+00$ | $1.99 \mathrm{e}+00$ |
| $250.0-275.0$ | $2.00 \mathrm{e}+00$ | $1.94 \mathrm{e}+00$ | $2.06 \mathrm{e}+00$ |
| $275.0-300.0$ | $1.89 \mathrm{e}+00$ | $1.83 \mathrm{e}+00$ | $1.95 \mathrm{e}+00$ |
| $300.0-325.0$ | $1.82 \mathrm{e}+00$ | $1.76 \mathrm{e}+00$ | $1.88 \mathrm{e}+00$ |
| $325.0-350.0$ | $1.82 \mathrm{e}+00$ | $1.77 \mathrm{e}+00$ | $1.88 \mathrm{e}+00$ |
| $350.0-375.0$ | $1.90 \mathrm{e}+00$ | $1.84 \mathrm{e}+00$ | $1.96 \mathrm{e}+00$ |

# Determination of Z, A, and Initial Energy from Voyager <br> HET Events <br> Appendix A --- Voyager 1 A-Stopping 

Boldface characters are computer output and Italics are comments. Standard Roman is your input.
cd/usr/erc/voyager/box
vrebox
Begin and end volumes : 1900000 ascii control file name : vihethas.box
input? 1
output? 2
What is name of output tape? <cr>
input? -1
In'ed to/usr/voyager/crspro/box/vrebox
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
cd/usr/voyager/crspro/ecal
ecall_lowz
Begin and end volumes : 1900000
input? 1
output? 2
What is name of output tape? <cr> input? -1
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
cd/usr/erc/voyager/astop
gainball vlgainal. 39
Begin and end volumes : 1900000
input? 1
Use appropriate or most recent gain file
for full tape
Input Tape Drive number
<input?>-1
Will only ask for second input on SHORT tapes
Pick up the two plots generated by gainball and label them with the CRS tape number and the gain file used. Match the plots with the theoretical track plots in the file folder "HET A-Stopping $H \& H e$ Final". If the tracks line up rename the gain file to reflect the additional CRS tape (vigaina1.s9 $=>$ v1gaina1.40). If there is not a match, op the gain file into a new file (cp vigaina1.99 v1gaina40), and adjust the gains in the new file to get the tracks to match. THEN RERUN GAINBAL1. If it still doesn't work, then you'll have to do it again.
eincalas v1gainal. 40
Begin and end volumes : 1900000
ascii control file name : v1h2a.ein
input? 1
output? 2
What is name of output tape? <cr> input? -1
1804 H, 32 H2, 47 Helium-3, 997 Helium-4

Note new gain file name
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
Gives number of good events
runfluxa
Number of time periods to sum : 2 The use of multiple time periods allows you to
Beginning and end volumes for each time period :

405601405792
405805410592
ascii control file name : vihas.flux
input? 1
<input?>-1 This will depend upon how many tapes are involved
What is the output file name for Hydrogen? v1has.h.88e
What is the output file name for Deuterium? v1has.h2.88e
What is the output file name for Helium-3? v1has.he3.88e
What is the output file name for Helium-4? v1has.he4.88e

The program writes one output file for each of masses in the vihas.flux file. For more information on the v1has. flux file, how to use these output files, and the naming conventions for the output file, see the section on Flux Calculation pages 11-13.

# Determination of Z, A, and Initial Energy from Voyager HET Events 

Appendix B --- Voyager 1 B-Stopping

Boldface characters are computer output and Italics are comments. Standard Roman is your input.
cd/usr/erc/voyager/box
vrebox
Begin and end volumes: 1900000
ascii control file name : vihethbs.box
imput? 1
output? 2
What is name of output tape? <cr>
input?-1
cd/usr/voyager/crspro/ecal
ecall_lowz
Begin and end volumes : 1900000
input? 1
output? 2
What is name of output tape? <cr> input?-1
cd /usr/erc/voyager/bstop
gainball vlgainb32.39
Begin and end volumes : 1900000
input? 1
<input?>-1
for full tape
In'ed to /usr/voyager/crspro/box/vrebox
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input

Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input

Use appropriate or most recent gain file for full tape
Input Tape Drive number
Will only ask for second input on SHORT tapes

Pick up the two plots generated by gainball and label them with the CRS tape number and the gain file used. Match the plots with the theoretical track plots in the file folder "HET B-Stopping H\& He Final". If the tracks line up rename the gain file to reflect the additional $C R S$ tape (vigainb32.39 $=>$ v1gainbs2.40). If there is not a match, cp the gain file into a new file (cp v1gainbs2.99 v1gainb40), and adjust the gains in the new file to get the tracks to match. THEN RERUN GAINBAL1. If it still doesn't work, then you'll have to do it again.
eincalbs v1gainb32.40
Begin snd end volumes : 1900000
ascii control file name: v1h2b.ein
input? 1
output? 2
What is name of output tape? <cr>
input? -1
1804 H, 32 H2, 47 Helium-3, 997 Helium-4

Note new gain file name
for full tape
Input Tape Drive number Output Tape Drive number
Never used
Ends chapter-verse input
Gives number of good events
runfluxb
Number of time periods to sum : 2 The use of multiple time periods allows you to Beginning and end volumes for each time period :
405601405792 avoid problems such as rate spikes
405805410592
ascii control file name : v1hbs.flux See Appendix $G$ for file format
input? 1
<input?>-1 This will depend upon how many tapes are involved
What is the output file name for Hydrogen? v1hbs.h.88e
What is the output file name for Deuterium? vlhbs.h2.88e
What is the output file name for Helium-3? vihbs.he3.88e
What is the output file name for Helium-4? v1hbs.he4.88e

The program writes one output file for each of masses in the vihbs.flux file. For more information on the vihbs.flux file, how to use these output files, and the naming conventions for the output file, see the section on Flux Calculation, pages 11-13.

# Determination of Z, A, and Initial Energy from Voyager HET Events <br> Appendix C --- Voyager 1 Penetrating 

Boldface characters are computer output and Italics are comments. Standard Roman is your input.
cd/usr/erc/voyager/box
vrebox
In'ed to /usr/voyager/crspro/box/vrebox
Begin and end volumes : 1900000
ascii control file name : vlhetpen.box
input? 1
output? 2
What is name of output tape? <cr> imput?-1
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
cd/usr/voyager/crspro/ecal
ecal1_lowz
Begin and end volumes: 1900000
input? 1
output? 2
What is name of output tape? <cr>
input? -1
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
cd/usr/erc/voyager/pen
gainball v1gainp1.39
Begin and end volumes : 1900000
input? 1
<input?>-1

> Use appropriate or most recent gain file
> for full tape
> Input Tape Drive number
> Will only ask for second input on SHORT tapes

Pick up the two plots generated by gainball and label them with the CRS tape number and the gain file used. Match the plots with the theoretical track plots in the file folder "HET Penetrating H 8 He Final". If the tracks line up rename the gain file to reflect the additional CRS tape (vigainp1.99 => v1gainp1.40). If there is not a match, cp the gain file into a new file (cp v1gainp1.39 v1gainp40), and adjust the gains in the new file to get the tracks to match. THEN RERUN GAINBAL1. If it still doesn't work, then you'll have to do it again.
eincalpen v1gainp1.40
Begin and end volumes : 1900000
ascii control file name : vih2p.ein input? 1
output? 2
What is name of output tape? <cr> input? -1
1804 H, 997 Helium Analyzed

Note new gain file name
for full tape

Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
Gives number of good events
runfluxp
Number of time periods to sum : 2
The use of multiple time periods allows you to
Beginning and end volumes for each time period :
405601405792
avoid problems such as rate spikes
405805410592
ascii control file name : v1hpn.flux
input? 1
See Appendix G for file format
<input?>-1 This will depend upon how many tapes are involved
What is the output file name for Hydrogen? v1hpn.h. 88 e
What is the output file name for Helium? v1hpn.he4.88e

Then you have to do runfluxp again to get the low energy point (79-87 MeV) separated into $A$ and $B$ directions.
runfluxp
Number of time periods to sum : 2
The use of multiple time periods allows you to
Beginning and end volumes for each time period :
405601405792 avoid problems such as rate spikes
405805410592
ascii control file name : v1low.flux See Appendix G for file format
input? 1
<input?>-1
Input Tape Drive number

What is the output file name for Hydrogen-A? vlhpna.h.88e
What is the output file name for Hydrogen-B? vlhpnb.h.88e

The program writes one output file for each of masses in the flux file. For more information on the .flux file, how to use these output files, and the naming conventions for the output file, see the section on Flux Calculation, pages 11-13.

# Determination of Z, A, and Initial Energy from Voyager HET Events <br> Appendix D --- Voyager 2 A-Stopping 

Boldface characters are computer output and Italics are comments. Standard Roman is your input.
cd/usr/erc/voyager/box
vrebox
Begin and end volumes: 1900000
ascii control file name : v2hethas.box
imput? 1
output? 2
What is name of output tape? <cr>
input?-1
cd /usr/voyager/crspro/ecal
ecal2_lowz
Begin and end volumes : 1900000
input? 1
output? 2
What is name of output tape? <cr> input? -1
cd /usr/erc/voyager/astop
gainbal2 v2gaina24.43
Begin and end volumes : 1900000
input? 1
<input?>-1

In'ed to /usr/voyager/crspro/box/vrebox
for full tape

Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input

Use appropriate or most recent gain file
for full tape
Input Tape Drive number Will only ask for second input on SHORT tapes

Pick up the two plots generated by gainbal2 and label them with the CRS tape number and the gain file used. Match the plots with the theoretical track plots in the file folder "HET A-Stopping $H \& H e$ Final". If the tracks line up rename the gain file to reflect the additional CRS tape (v2gaina24.43 => v2gaina24.44). If there is not a match, cp the gain file into a new file (cp v2gaina24.43 v2gaina44), and adjust the gains in the new file to get the tracks to maich. THEN RERUN GAINBAL2. If it still doesn't work, then you'll have to do it again.
eincalas v2gaina24.44
Begin and end volumes : 1900000
ascii control file name : v2hla.ein
input? 1
output? 2
What is name of output tape? $-<c r>$ input? -1
1804 H, 32 H2, 47 Helium-3, 997 Helium-4

Note new gain file name
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
Gives number of good events
runfluxa
Number of time periods to sum : 2 The use of multiple time periods allows you to
Beginning and end volumes for each time period :
405601405792
avoid problems auch as rate spikes
405805410592
ascii control file name : v2has.flux See Appendix $G$ for file format
input? 1
Input Tape Drive number
<input?>-1 This will depend upon how many tapes are involved
| What is the output file name for Hydrogen? v2has.h.88e
What is the output file name for Deuterium? v2has.h2.88e
What is the output file name for Helium-3? v2has.he3.88e
What is the output file name for Helium-4? v2has.he4.88e

The program writes one output file for each of masses in the v2has.flux file. For more information on the v2has.flux file, how to use these output files, and the naming conventions for the output file, see the section on Flux Calculation, pages 11-13.

# Determination of Z, A, and Initial Energy from Voyager HET Events Appendix E --- Voyager 2 B-Stopping 

Boldface characters are computer output and Italics are comments. Standard Roman is your input.
cd /usr/erc/voyager/box
vrebox

## ln'ed to /usr/voyager/crspro/box/vrebox <br> for full tape

Begin and end volumes : 1900000
ascii control file name : v2hethbs.box
input? 1
output? 2
What is name of output tape? <cr>
input? -1
Input Tape Drive number
Output Tape Drive number
Never used
Ends ehapter-verse input
cd /usr/voyager/crspro/ecal
ecal2_lowz
Begin and end volumes : 1900000
input? 1
output? 2
What is name of output tape? <cr> input? -1
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
cd /usr/erc/voyager/bstop
gainbal2 v 2 gainb 28.43
Begin and end volumes : 1900000
input? 1
<input?>-1

> Use appropriate or most recent gain file
> for full tape
> Input Tape Drive number
> Will only ask for second input on SHORT tapes

Pick up the two plots generated by gainbal2 and label them with the CRS tape number and the gain file used. Match the plots with the theoretical track plots in the file folder "HET B-Stopping H \& He Final". If the tracks line up rename the gain file to reflect the additional CRS tape (v2gainb28.49 $=>$ v2gain628.48). If there is not a match, cp the gain fite into a new file (cp v2gain628.49 v2gainb44), and adjust the gains in the new file to get the tracks to match. THEN RERUN GAINBAL2. If it still doesn't work, then you'll have to do it again.
eincalbs v2gainb28.44
Begin and end volumes : 1900000 ascii control file name : v2hlb.ein input? 1
output? 2
What is name of output tape? <cr> input? -1
1804 H, 32 H2, 47 Helium-3, 997 Helium-4

Note new gain file name
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
Gives number of good events
runfluxb
Number of time periods to sum : 2 The use of multiple time periods allows you to
Beginning and end volumes for each time period:
405601405792
405805410592
: ascii control file name : v2hbs.flux
input? 1
avoid problems such as rate spikes
<input?>-1
See Appendix $G$ for file format
Input Tape Drive number
What is the output file name for Hydrogen? v2hbs.h. 88 e
What is the output file name for Deuterium? v2hbs.h 2.88 e
What is the output file name for Helium-3? v2hbs.he3.88e
What is the output file name for Helium-4? v2hbs.he4.88e
The program writes one output file for each of masses in the v2hbs.flux file. For more information on the v2hbs.flux file, how to use these output files, and the naming conventions for the output file, see the section on Flux Calculation, pages 11-13.

# Determination of Z, A, and Initial Energy from Voyager HET Events <br> Appendix F --- Voyager 2 Penetrating 

Boldface characters are computer output and Italics are comments. Standard Roman is your input.
cd /usr/erc/voyager/box
vrebox
Begin and end volumes : 1900000
ascii control file name : v2hetpen.box
input? 1
output? 2
What is name of output tape? <cr> input! -1
cd /usr/voyager/crspro/ecal
ecal2_lowz
Begin and end volumes : 1900000
input? 1
output? 2
What is name of output tape? <cr>
input? -1
cd /usr/erc/voyager/pen
gainbal2 v2gainp35.43
Begin and end volumes: 1900000
input? 1
<input?>-1

In'ed to /usr/voyager/crspro/box/vrebox for full tape

Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
for full tape
Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input

Use appropriate or most recent gain file
for full tape
Input Tape Drive number
Will only ask for second input on SHORT tapes

Pick up the two plots generated by gainbal2 and label them with the CRS tape number and the gain file used. Match the plots with the theoretical track plots in the file folder "HET Penetrating H 8 He Finaf". If the tracks line up rename the gain file to reflect the additional CRS tape (v2gainp95.48 => v2gainp35.44). If there is not a match, cp the gain file into a new file (cp v2gainp35.49 v2gainp44), and adjust the gains in the new file to get the tracks to match. THEN RERUN GAINBAL2. If it still doesn't work, then you'll have to do it again.
eincalpen v2gainp35.44
Begin and end volumes : 1900000
ascii control file name : v2h1p.ein input? 1
output? 2
What is name of output tape? <cr> input? -1
1804 H, 997 Helium Analyxed

Note new gain file name
for full tape

Input Tape Drive number
Output Tape Drive number
Never used
Ends chapter-verse input
Gives number of good events
runfluxp
Number of time periods to sum : 2
The use of multiple time periods allows you to
Beginning and end volumes for each time period :
405601405792
avoid problems such as rate spikes
405805410592
ascii control file name : v2hpn.flux See Appendix G for file format
input? 1
Input Tape Drive number
<input?>-1 This will depend upon how many tapes are involved
What is the output file name for Hydrogen? v2hpn.h.88e
What is the output file name for Helium? v2hpn.he4.88e

Then you have to do runfluxp again for the low energy (79-87 MeV) point to separate it into $A$ and $B$.
runfluxp
Number of time periods to sum : 2 The use of multiple time periods allows you to
Beginning and end volumes for each time period :
405601405792 avoid problems such as rate spikes
405805410592
ascii control file name : v2low.flux See Appendix $G$ for file format
input? 1
Input Tape Drive number
<input?>-1 This will depend upon how many tapes are involved
What is the output file name for Hydrogen-A? v2hpna.h.88e
What is the output file name for Hydrogen-B? v2hpnb.h.88e

The program writes one output file for each of masses in the .flux file. For more information on the .flux file, how to use these output files, and the naming conventions for the output file, see the section on Flux Calculation, pages 11-13.

